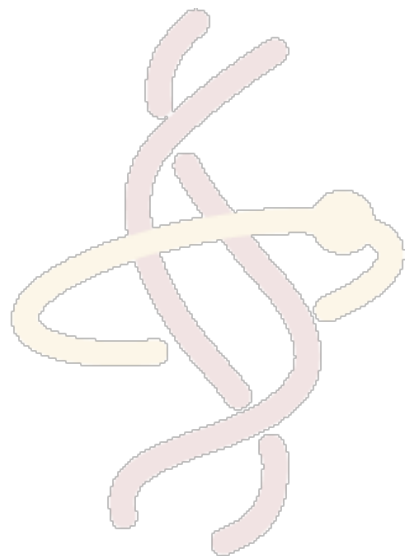


The laws of nature that we care about emerge through collective self-organization and really do not require knowledge of their component parts ... they owe their reliability to principles of organization rather than to microscopic rules.

- Robert Laughlin

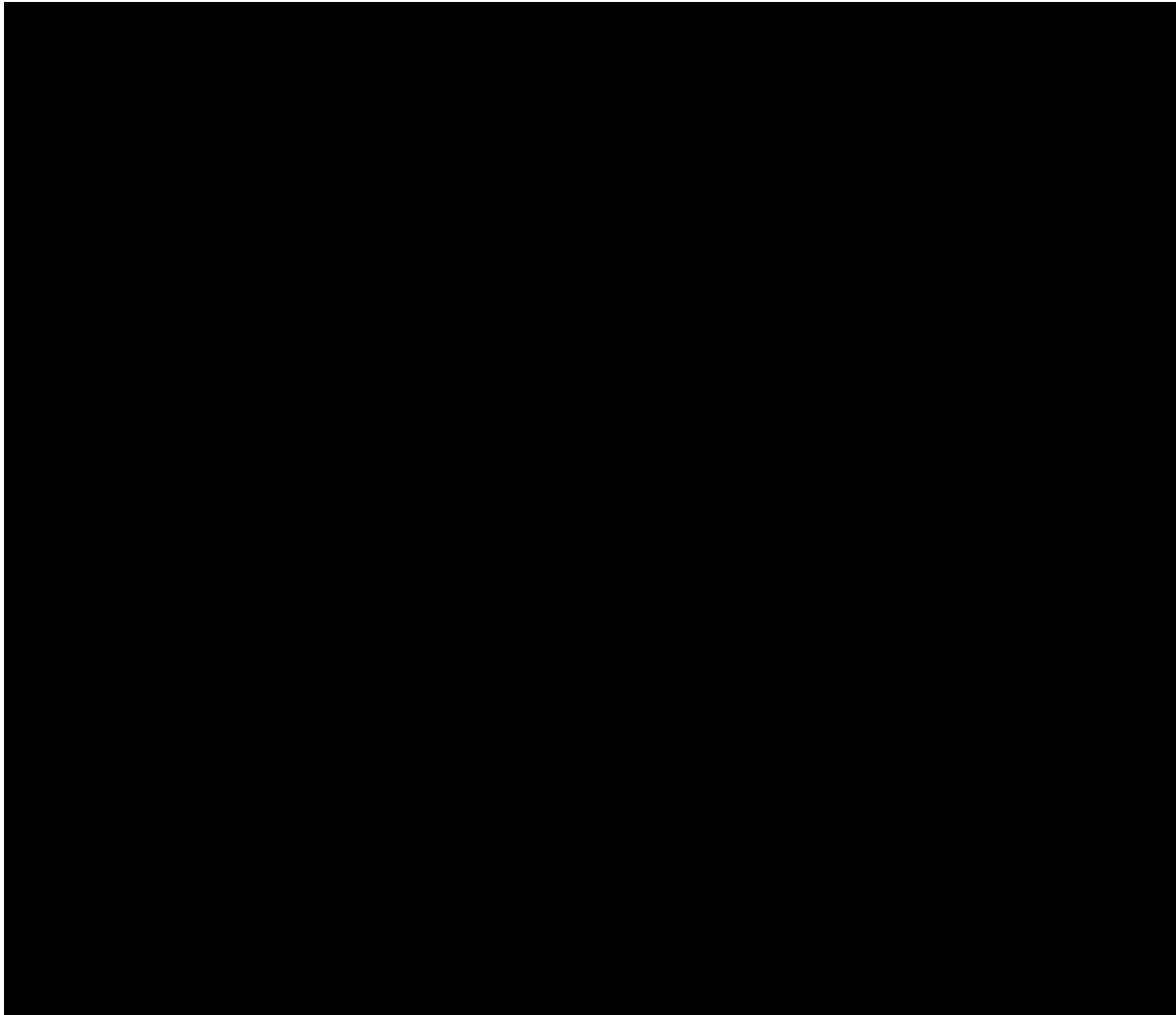
What can biology teach us about the condensed phase?

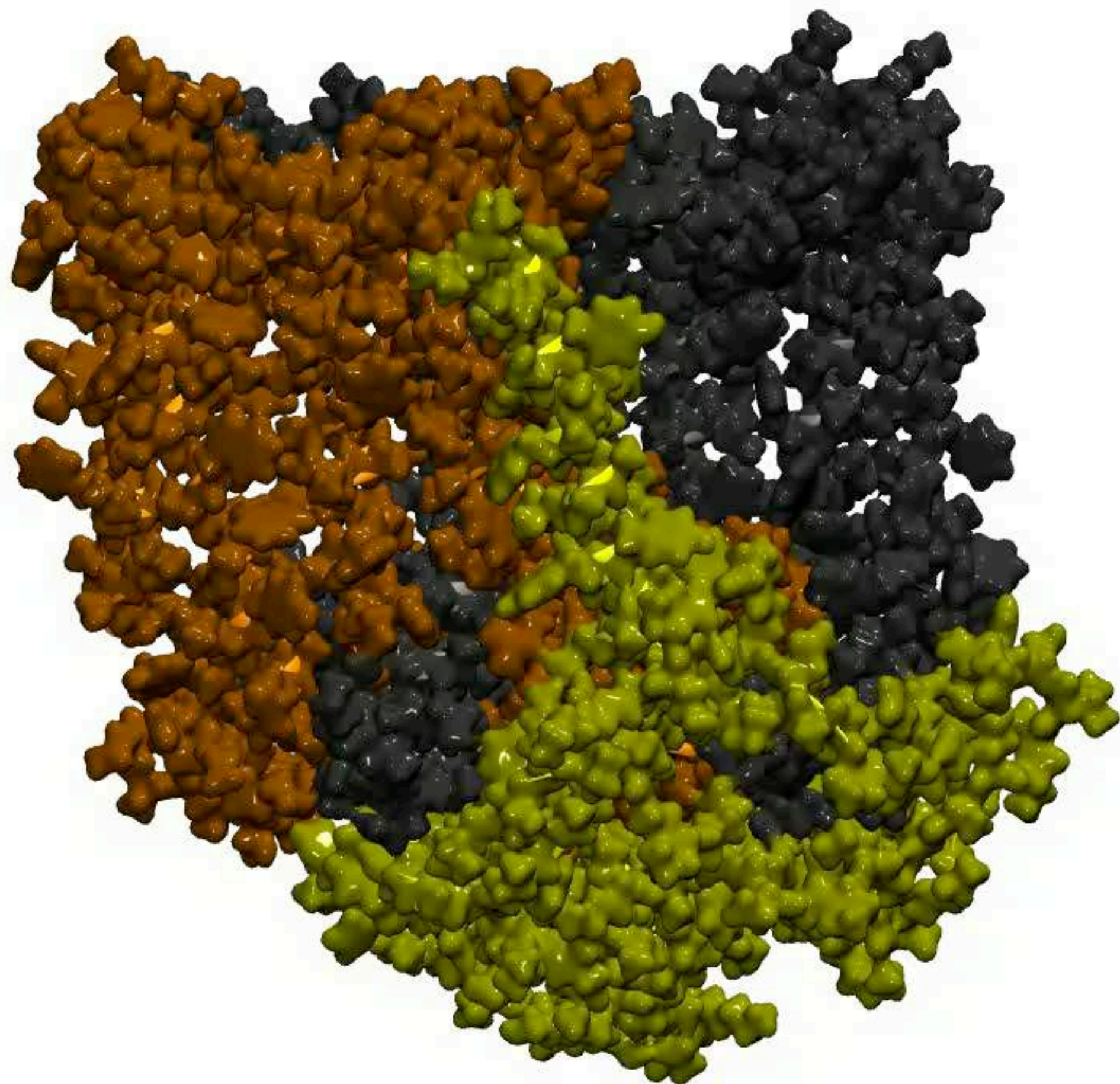
Tampa, Nov 13, 2014



Dmitry Matyushov
Center for Biological Physics
Department of Physics/Chemistry
Arizona State University

Bacterial reaction center: Are there certain general rules here?

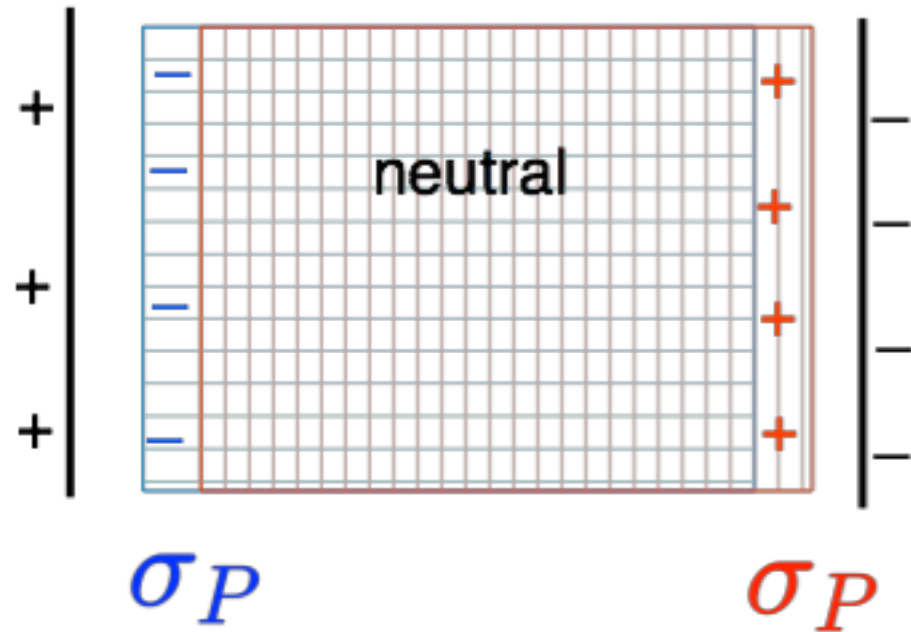




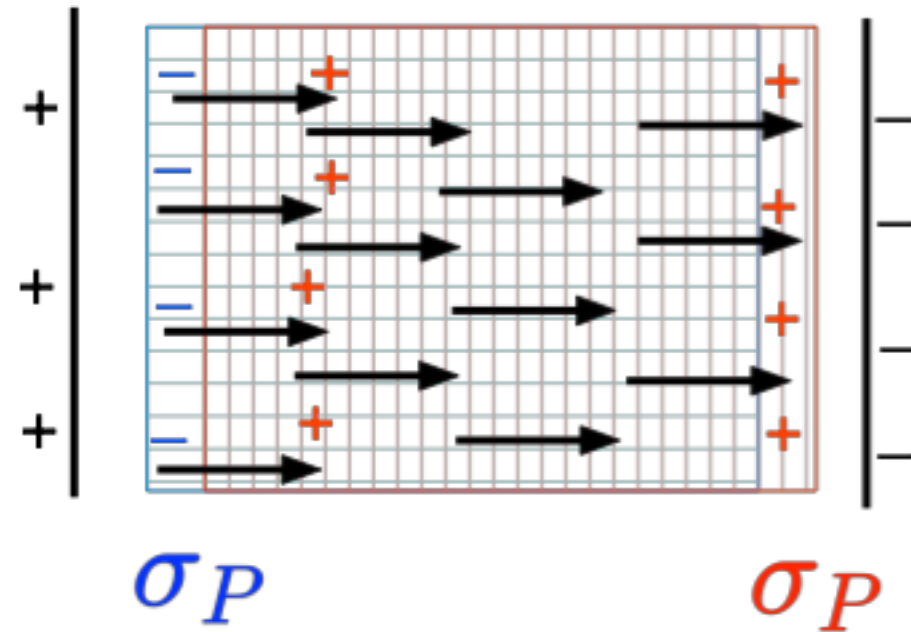
Electric Elasticity (Maxwell)

When an electromotive force acts on a dielectric, it puts every part of the dielectric into a polarized condition, in which its opposite sides are oppositely electrified.

“A dynamical Theory of Electromagnetic Field”, Maxwell, 1865.



Maxwell: Elastic deformation of positive vs negative liquid

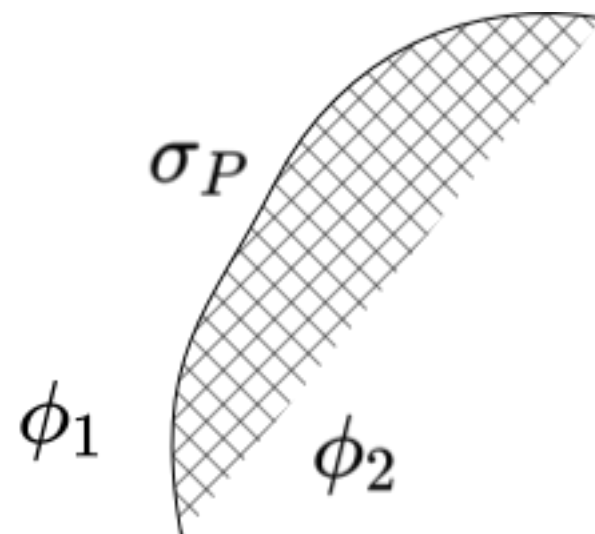


Debye: Picture of uniformly oriented dipoles

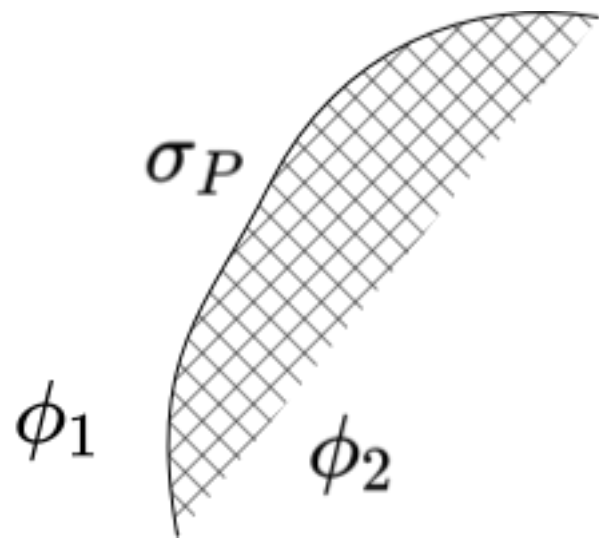
Boundary value problem:

$$\frac{\partial \Delta \phi}{\partial n} = -\Delta E_n = 4\pi \sigma_P$$

$$\Delta \phi = \phi_1 - \phi_2$$

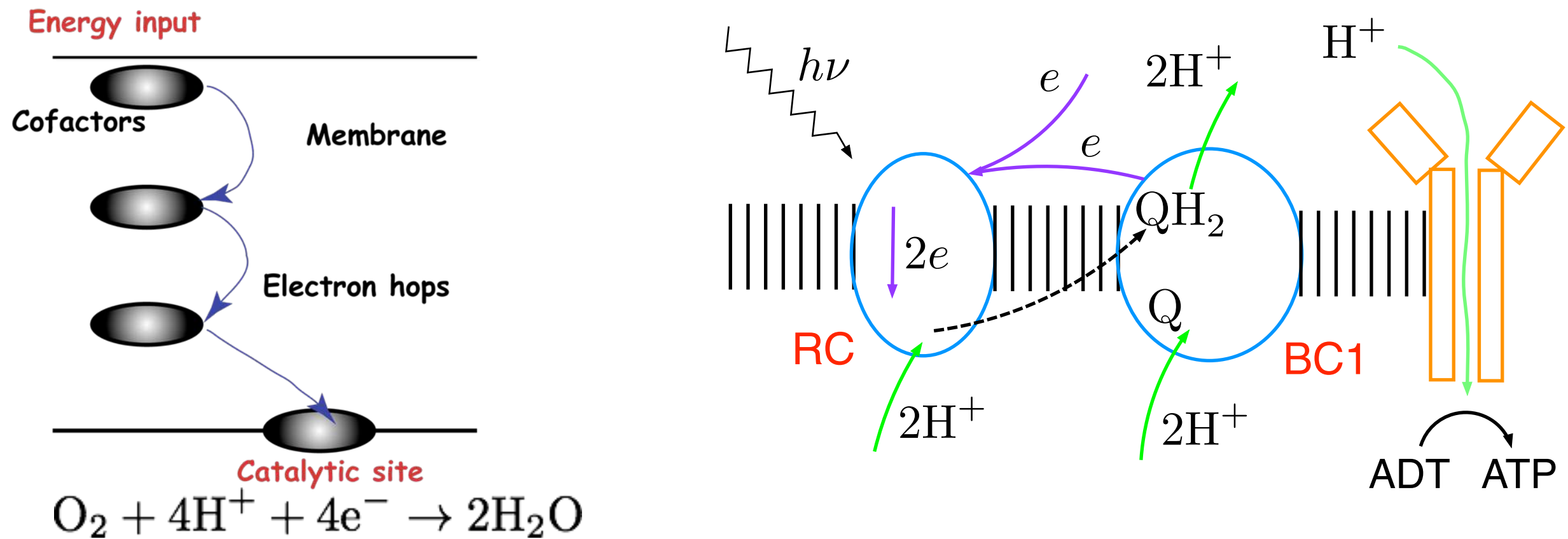


...the advancement of science depends on the discovery and development of exact ideas ... to warrant the deductions we may draw by the application of mathematical reasoning. - J. C. Maxwell



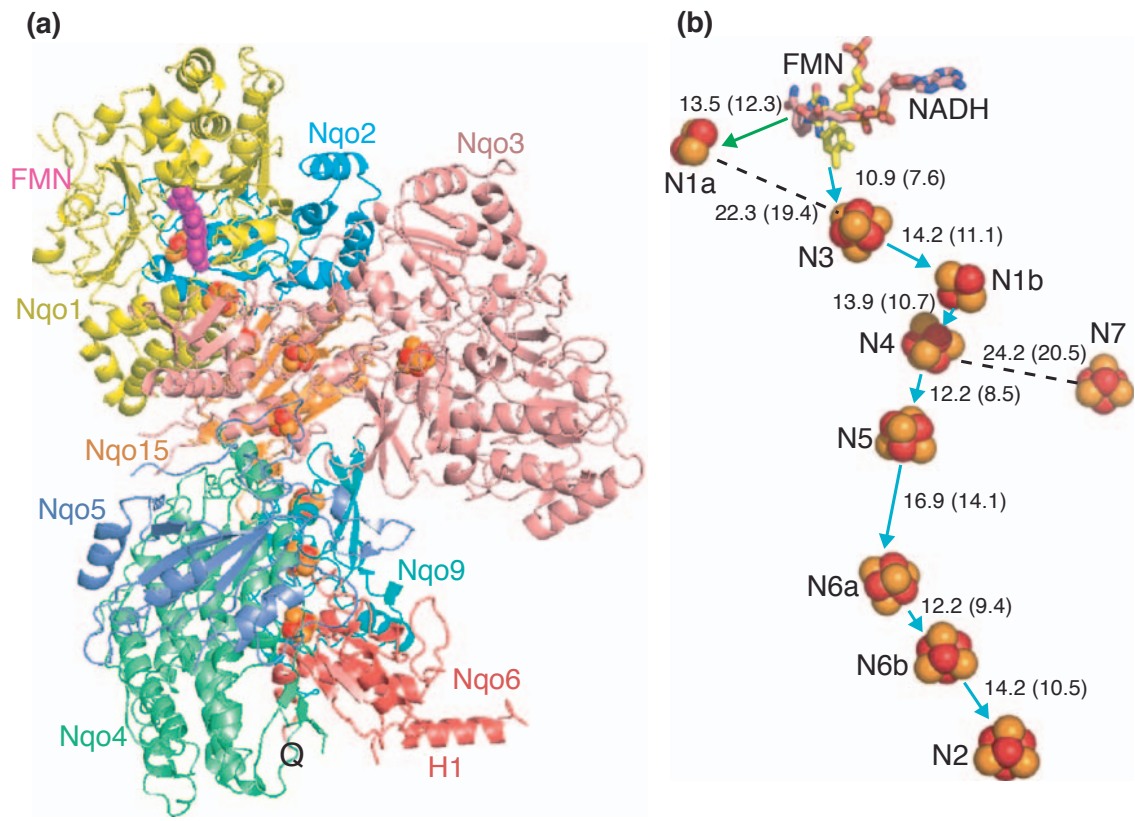
If the mathematics is universal, where do the specifics of the system come in?

Biology's energy chains

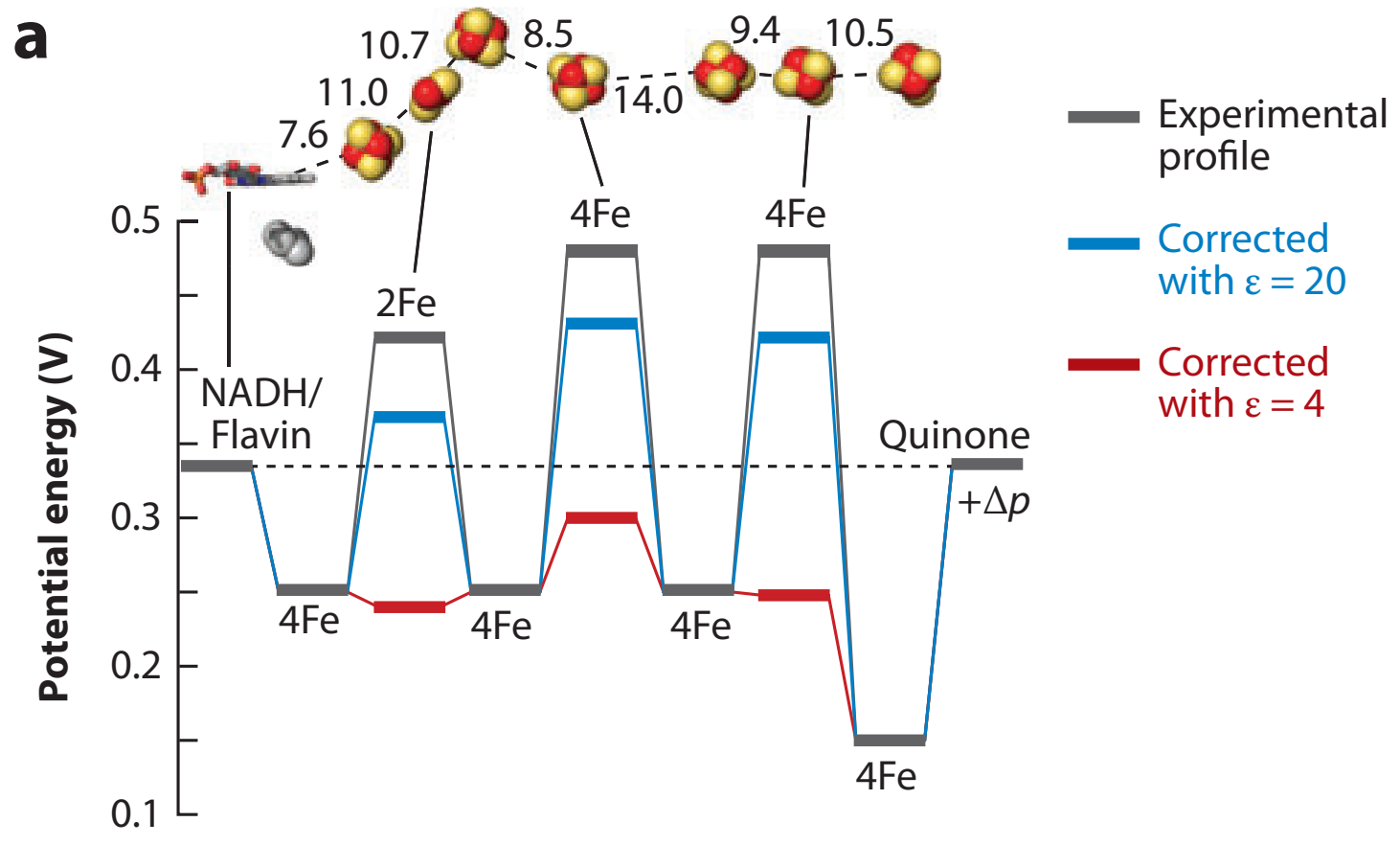


- 22 electron hops in mitochondria's membrane over the free energy span of 1.1 eV
- 8-9 electrons per one ATP produced
- ~25 kg of ATP produced daily in a human body

Mitochondria's Complex I

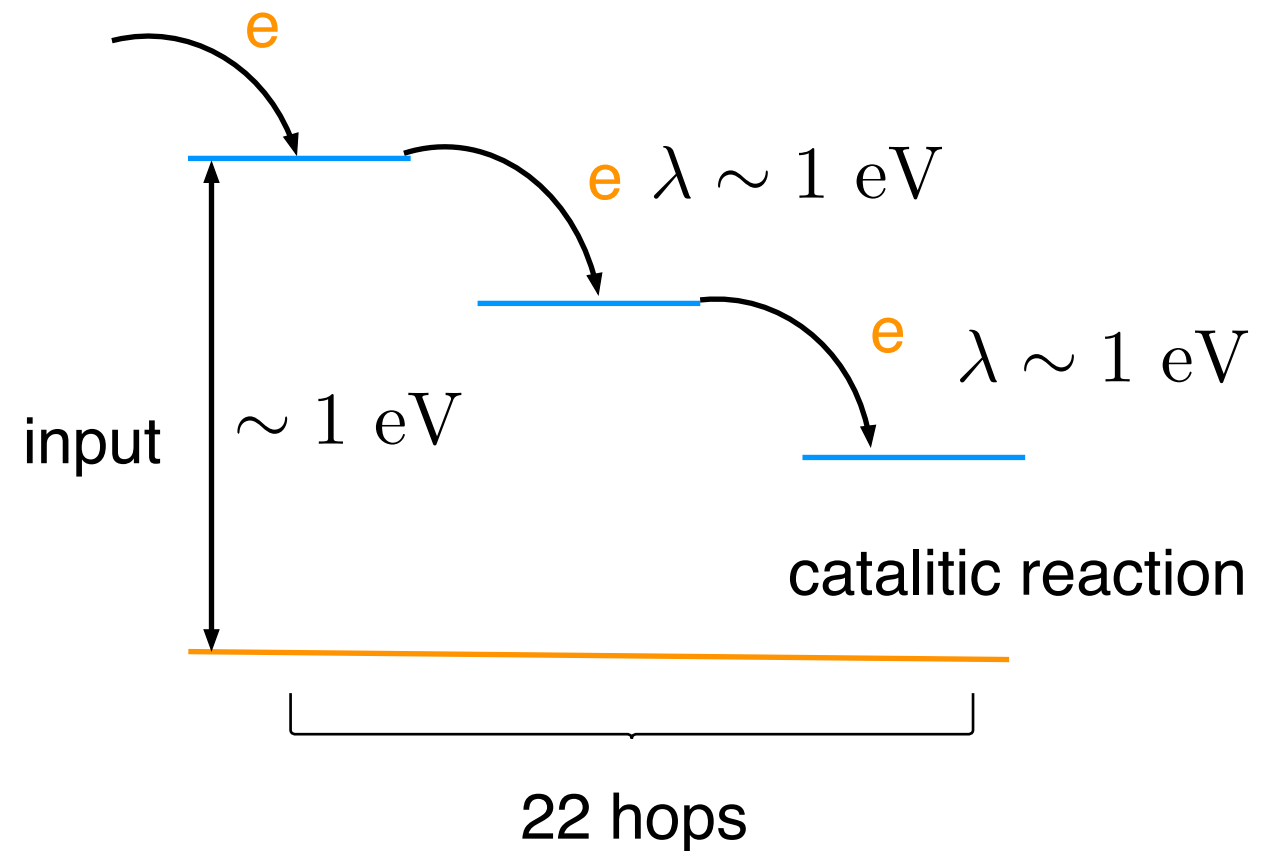
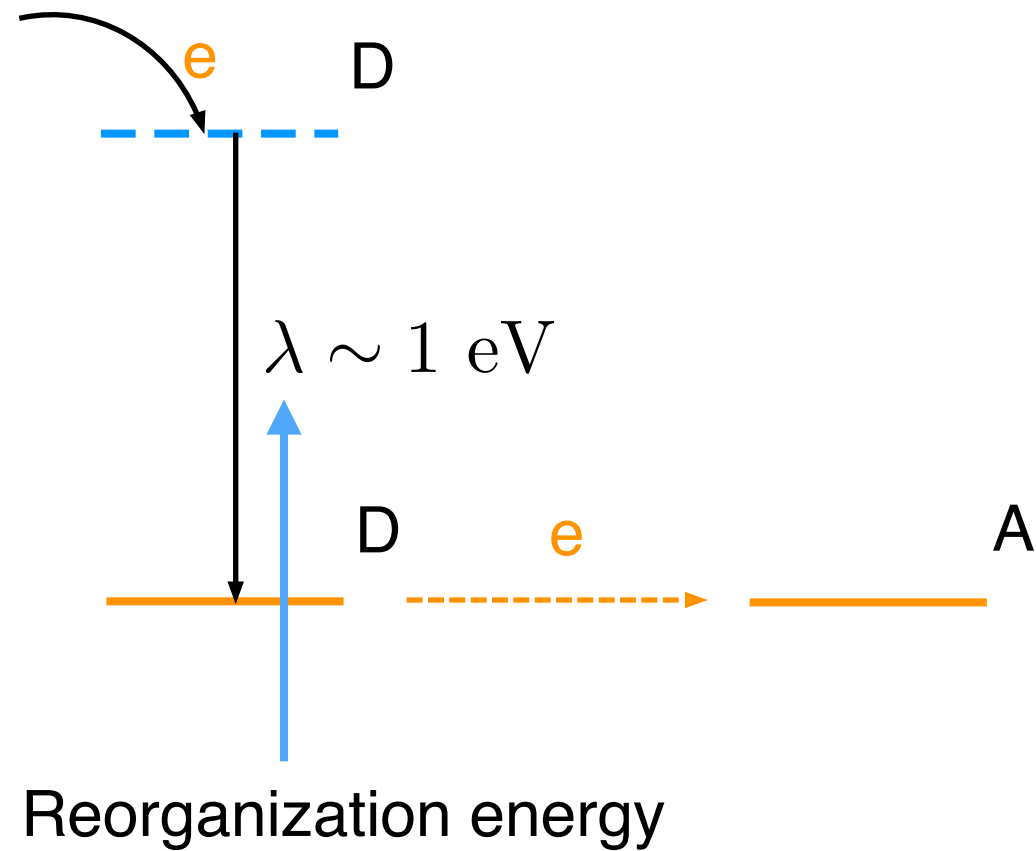


Efremov and Sazanov,
Curr. Opinion Struct. Biol. 2011, 21:532–540



J. Hurst, Annu. Rev. Biochem. 2013. 82:551–75

Energy balance

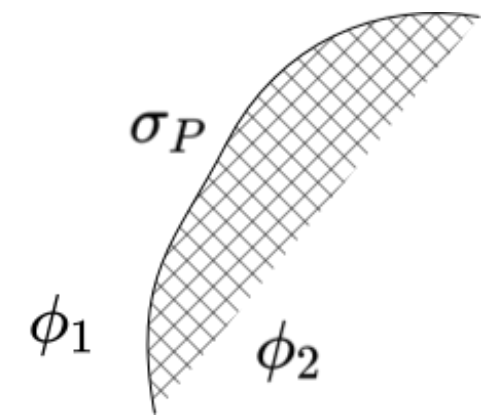


How does biology produce energy?

The mathematical framework seems to be OK, but we tend to put the “wrong” physics into it!

Is it thermodynamics only?

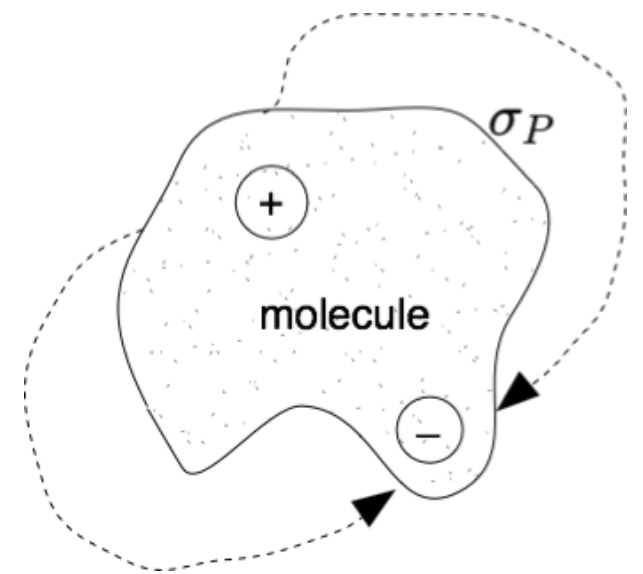
Can we view biological energy flow as a
canonical ensemble problem?



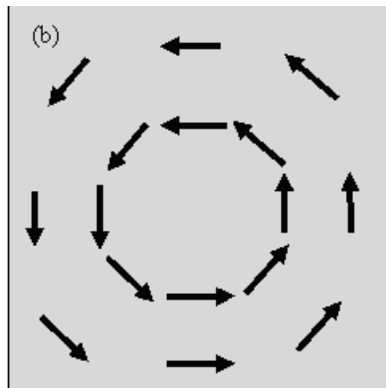
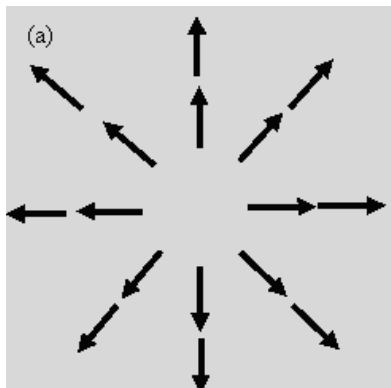
$$\frac{\partial \Delta \phi}{\partial n} = -\Delta E_n = 4\pi \sigma_P$$

Constitutive relation:

$$\mathbf{P} = \chi \mathbf{E}, \quad \underbrace{\mathbf{D}}_{\text{transverse}} = \epsilon \underbrace{\mathbf{E}}_{\text{longitudinal}}$$



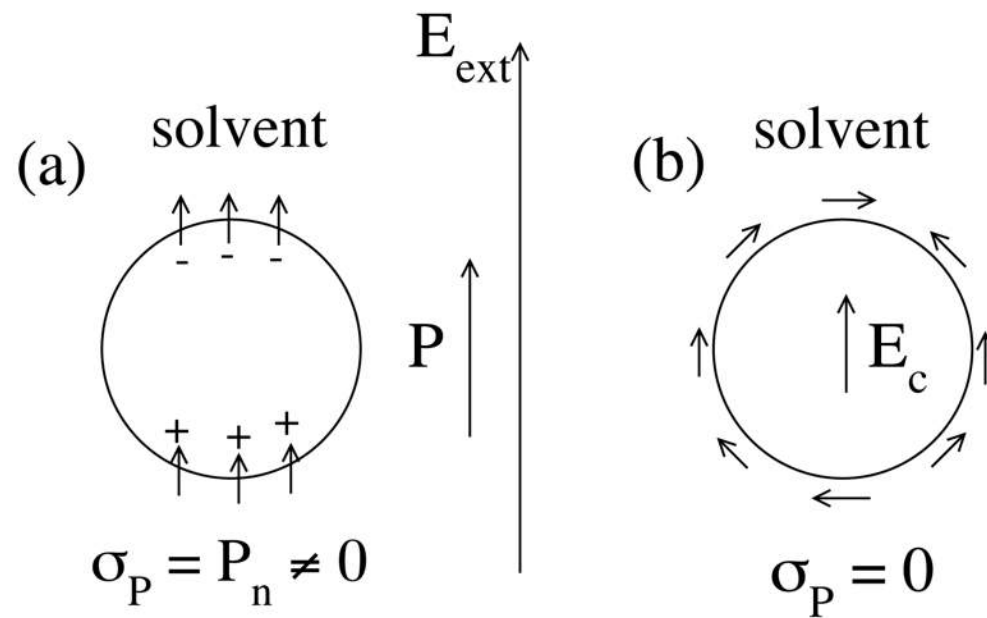
$$\nabla \cdot \mathbf{D} = 0$$



$$\nabla \times \mathbf{E} = 0$$

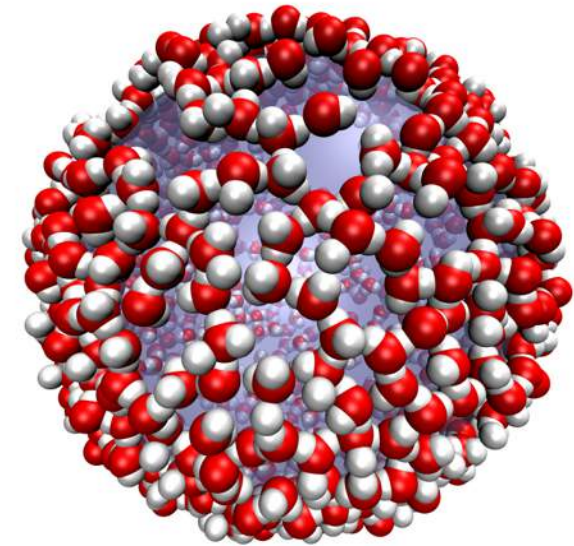
For non-uniform fields the constitutive relations connects fields of different symmetry

Cavity field



Kihara solute:

$$\phi_{0s}(r) = 4\epsilon_{0s} \left[\left(\frac{\sigma_{0s}}{r - R_{\text{HS}}} \right)^{12} - \left(\frac{\sigma_{0s}}{r - R_{\text{HS}}} \right)^6 \right]$$

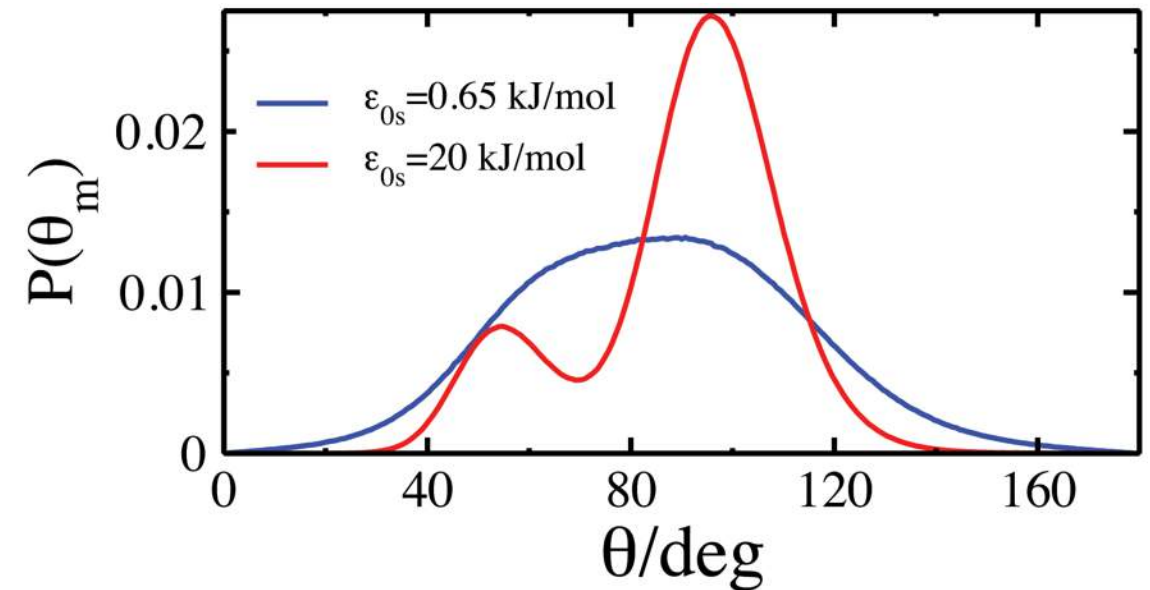


(a) Maxwell scenario:

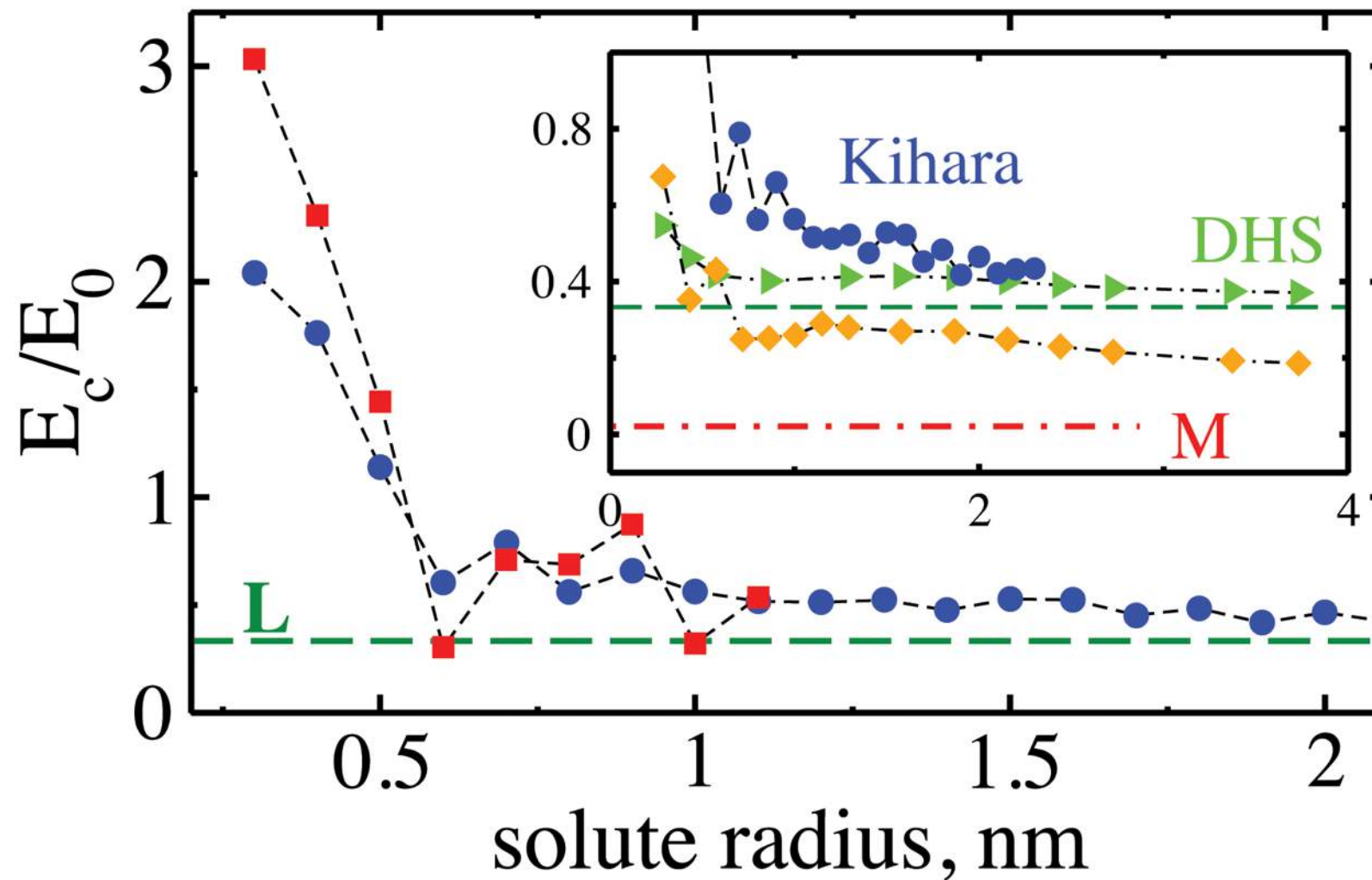
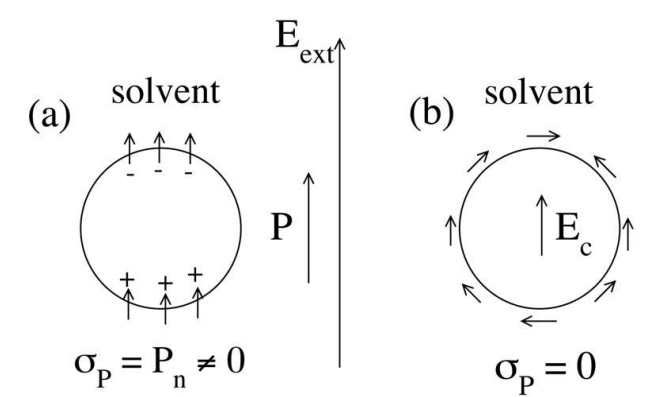
$$\frac{E_c}{E_{\text{ext}}} = \frac{3}{2\epsilon + 1}$$

(b) Lorentz scenario:

$$\frac{E_c}{E_{\text{ext}}} = \frac{\epsilon + 2}{3\epsilon}, \quad \sigma_P = 0$$



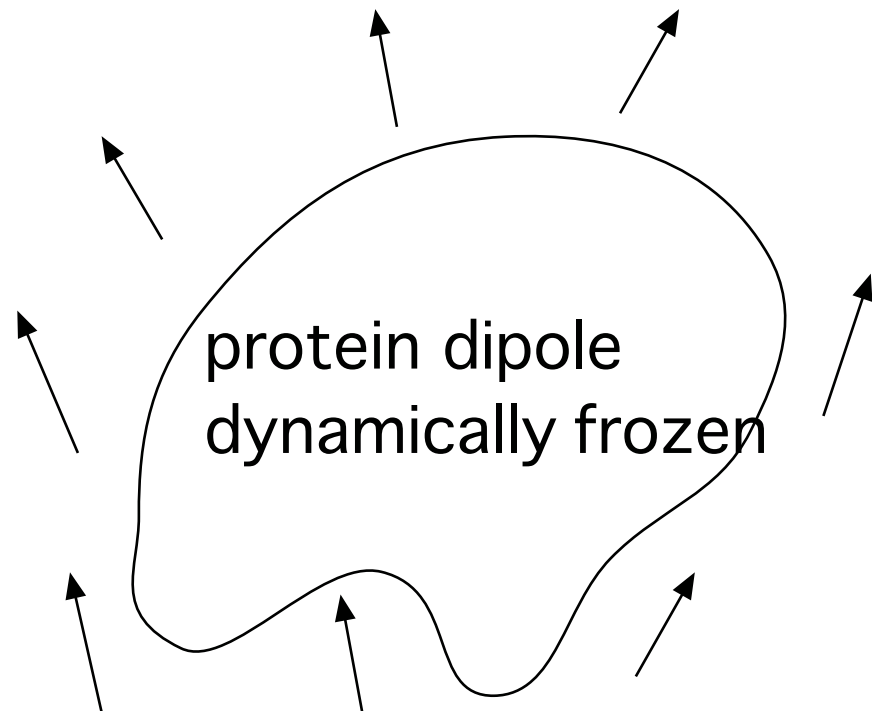
“Cavity” field inside a Kihara solute



Lorentz scenario (no surface polarization) is more consistent with the data

Dipole of the interface

Total induced dipole of the hydration layer



$$\alpha = M_0^{\text{int}} / M_0^{\text{M}}$$

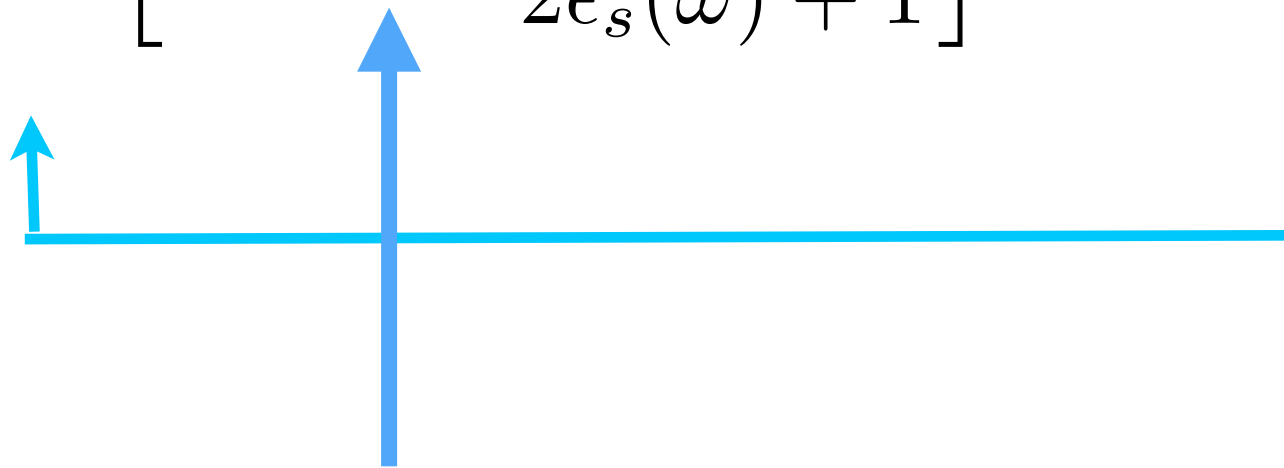
Parameter quantifying the deviation from the Maxwell scenario

Maxwell interface dipole projected on x-axis of the external field

High-frequency (THz) absorption

$$\alpha_{\text{abs}}(\omega) = \frac{4\pi\omega}{c} \frac{\chi''(\omega)}{\sqrt{1 + 4\pi\chi'(\omega)}}$$

$$\frac{\Delta\chi(\omega)}{\chi_s(\omega)} = -\eta_0 \left[1 + \alpha(\omega) \frac{\epsilon_s(\omega) - 1}{2\epsilon_s(\omega) + 1} \right]$$

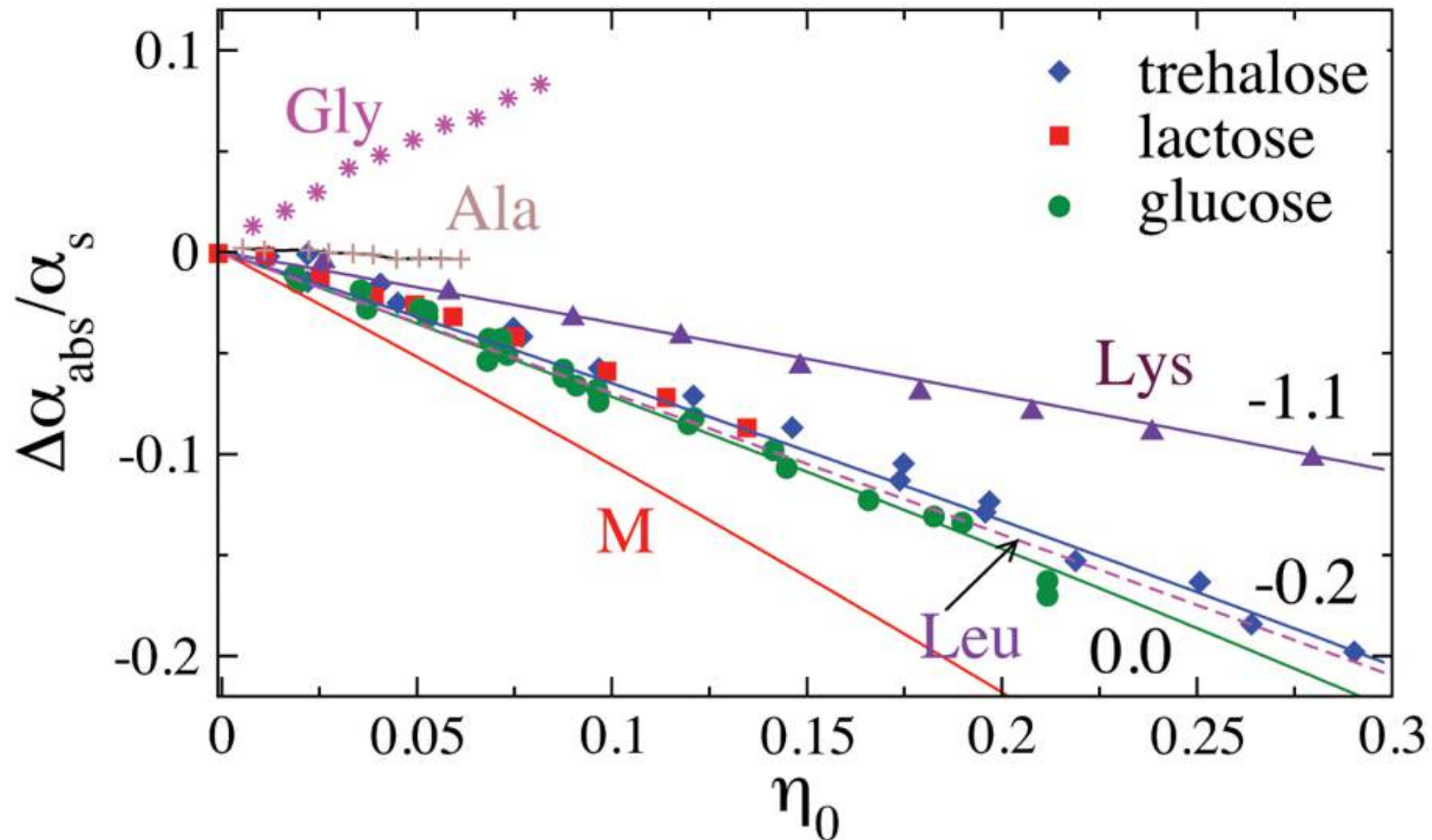


Deviation from the Maxwell prediction

Volume fraction of
solutes in solution

THz absorption of sugars and amino acids (aq)

Rotations of a large solute are dynamically frozen on the THz time-scale, solutes are approximated by dielectric voids



Heyden et al, JACS 130 (2008) 5773

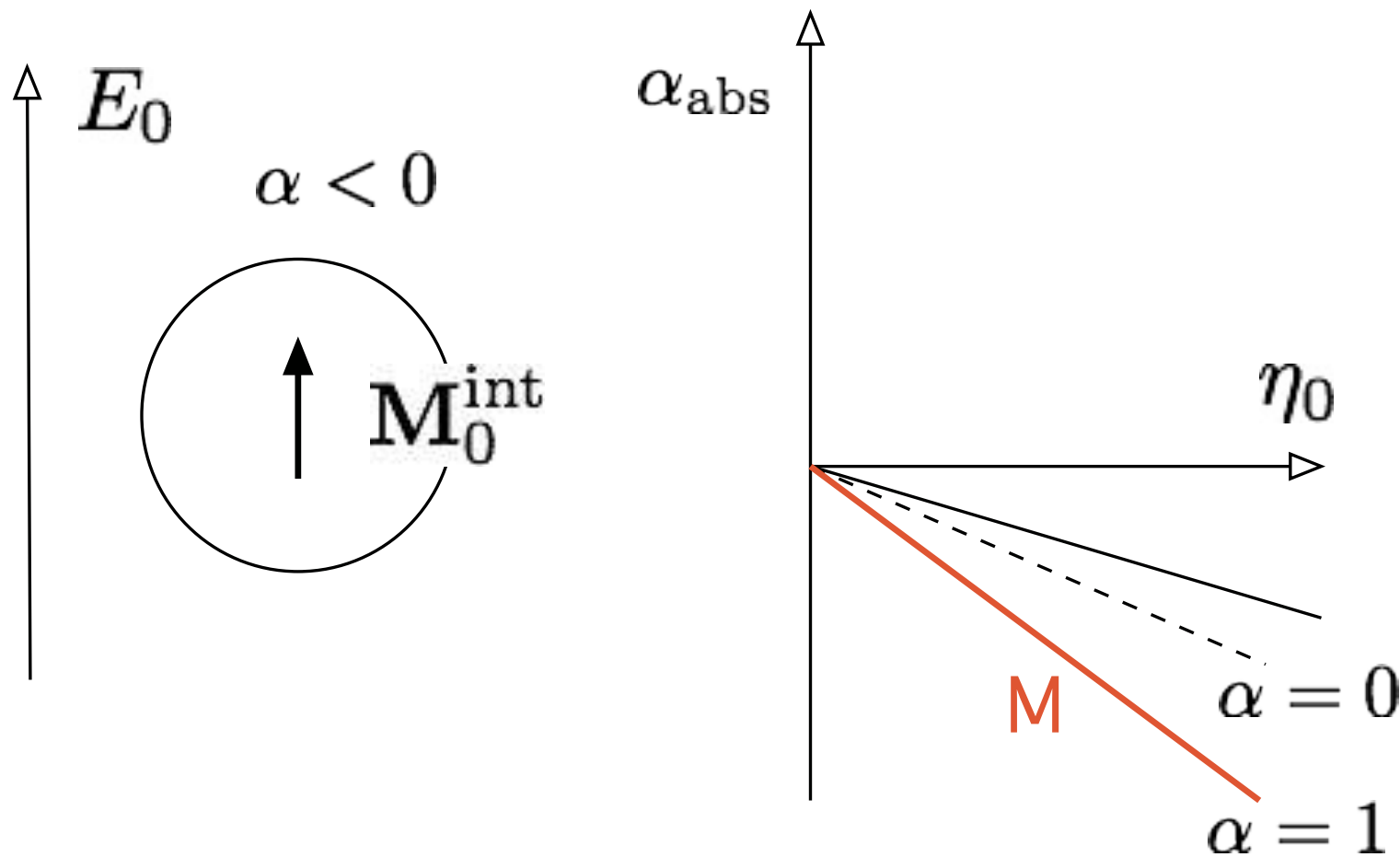
Niehues et al. Farad. Disc. Chem. Soc. 150 (2011) 193

What does THz absorption tell us?

Maxwell: $\alpha = 1$

Sugars: $\alpha = -0.2 - 0$, **no interface dipole!**

Amino acids: $\alpha = -5 - -0.1$, **interface dipole opposite to the field!**



Weakly **hydrophilic** solutes (glucose) leave no footprint, **hydrophilic** solute enhance the polarity relative to water.

Lysozyme: cavity field

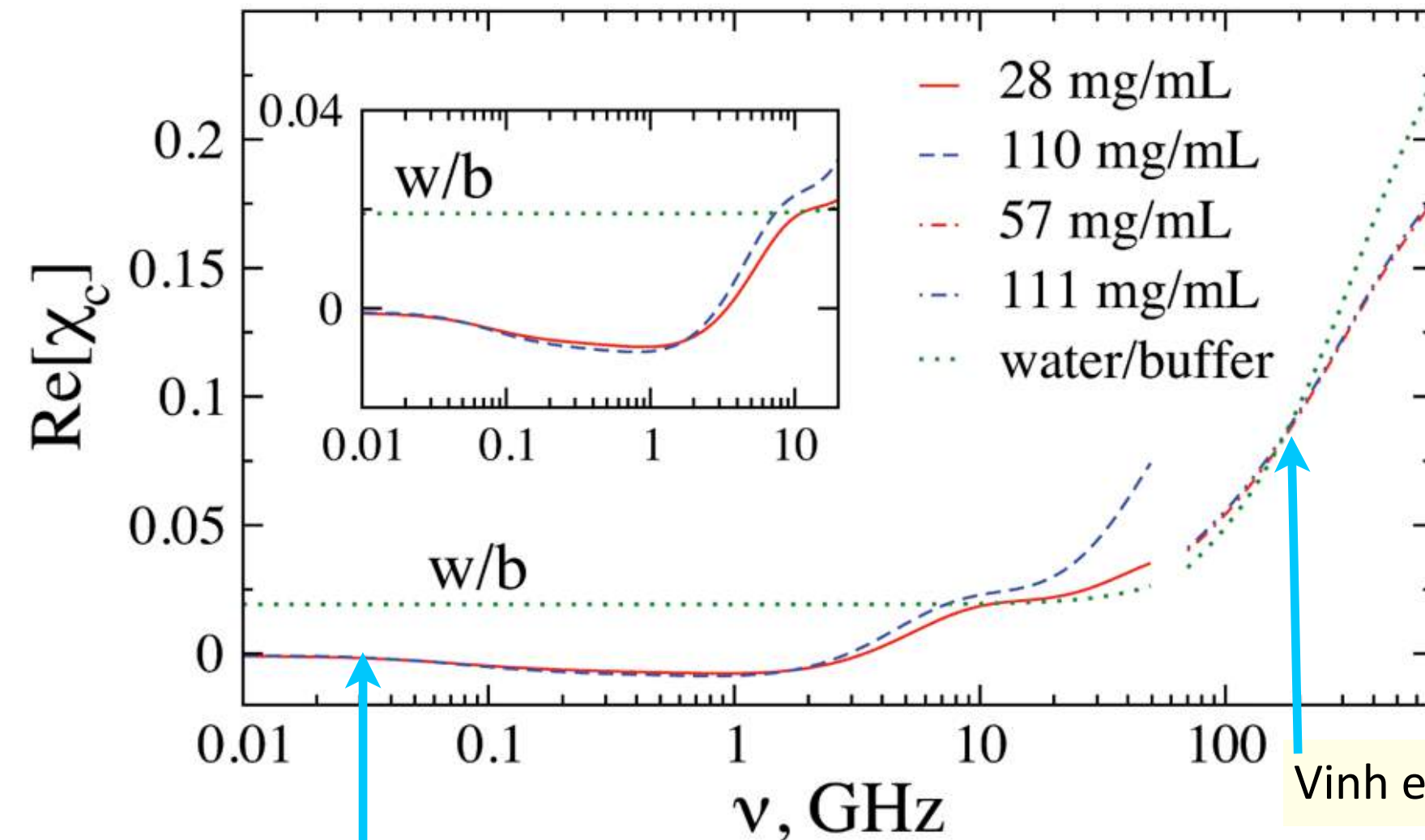
$$\frac{\epsilon_s(\omega)}{\epsilon_{\text{mix}}(\omega)} = 1 - 3\eta_0 + 3\eta_0\epsilon_s(\omega)\chi_c(\omega)(1 - y_0(\omega))$$

Experimental input

Volume fraction

Response of the lysozyme dipole, taken from MD

Cavity field response



$$\chi_c^M = \frac{3}{2\epsilon_s + 1}$$

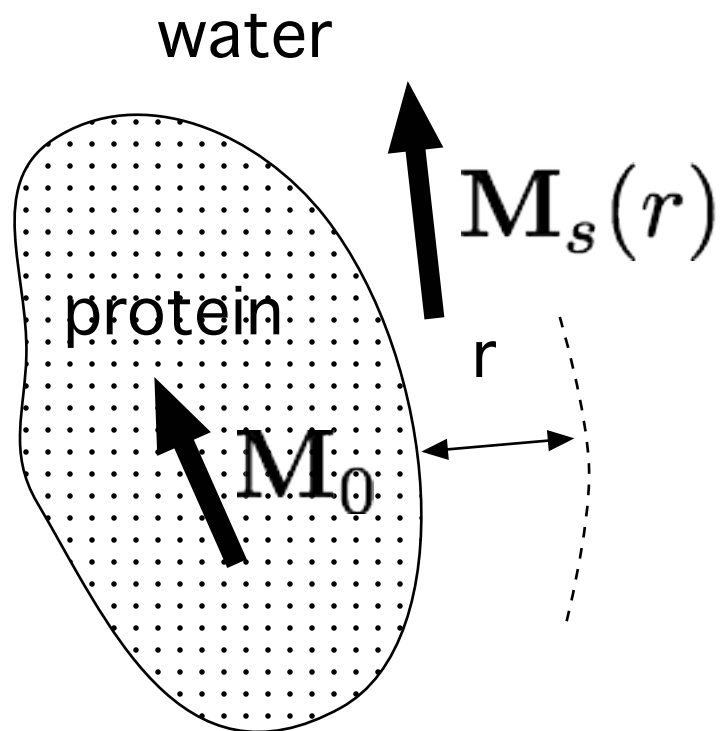
Maxwell scenario

Cametti et al., JPCB 115 (2011) 7144

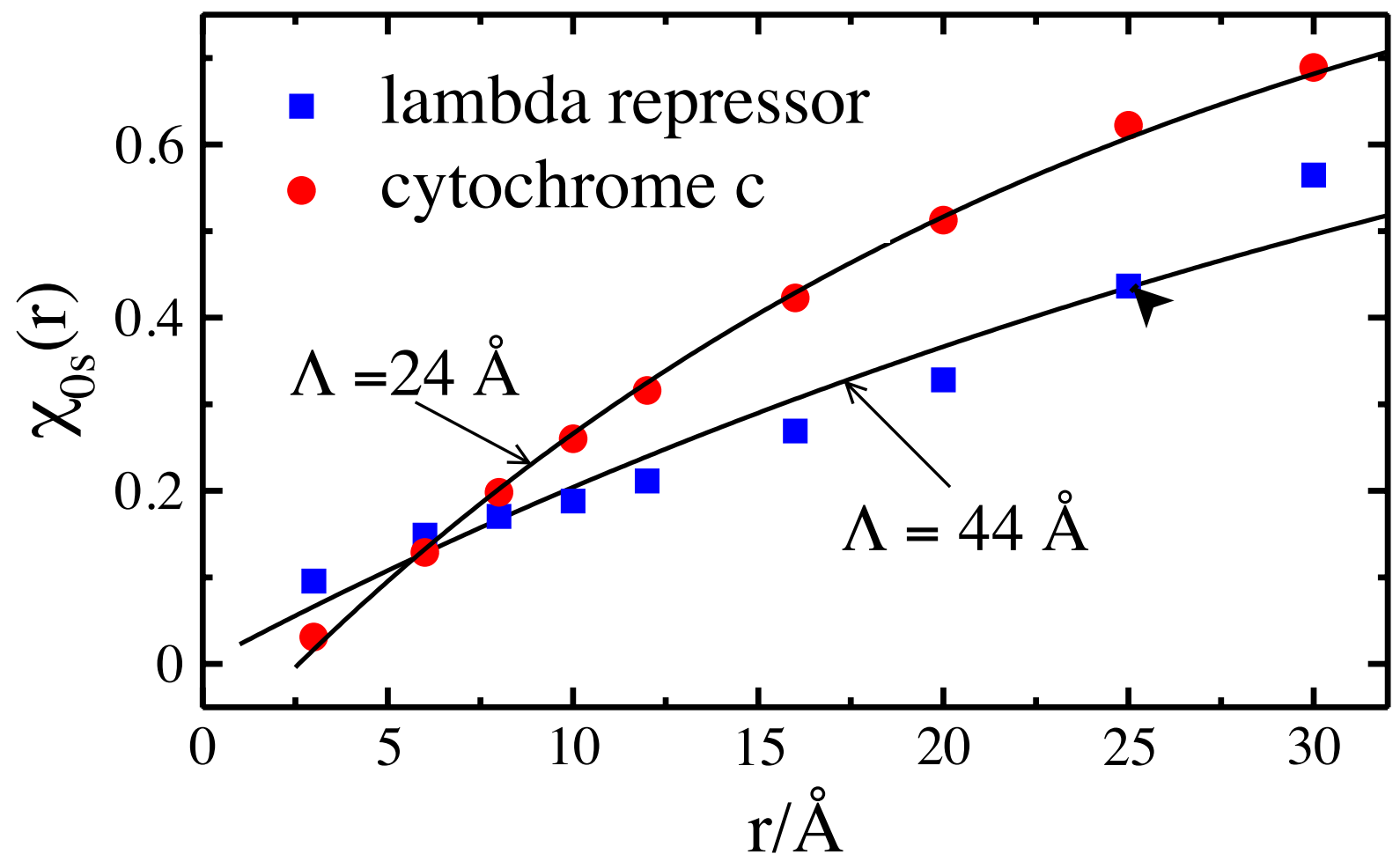
Vinh et al, JACS 133 (2011) 8942

DVM, JPCM 24 (2012) 325105

Range of protein-water correlations

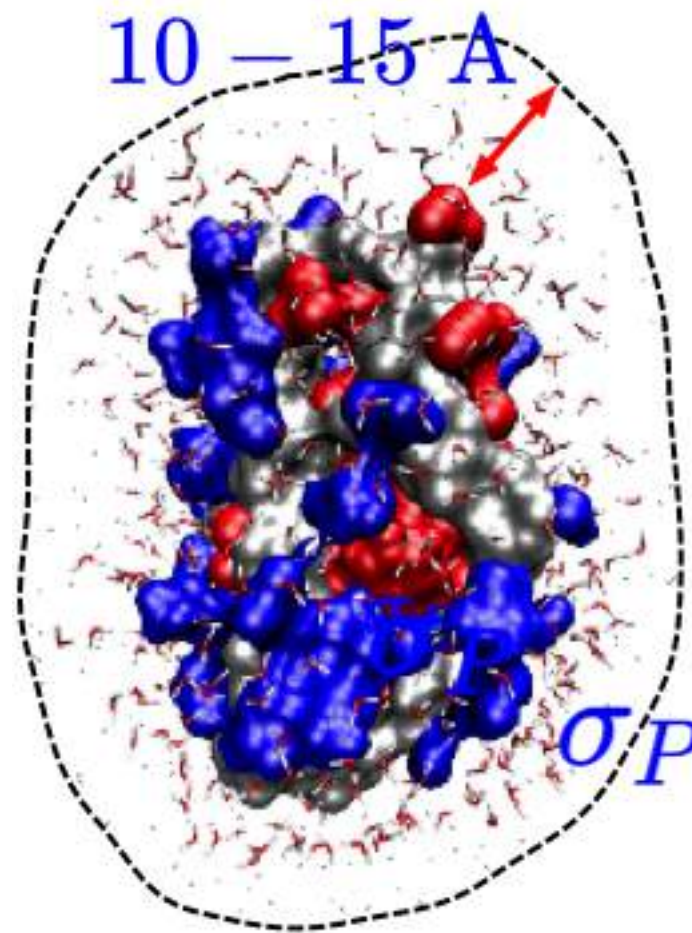


$$\chi_{0s}(r) \propto \langle \delta \mathbf{M}_0 \cdot \delta \mathbf{M}_s(r) \rangle$$



$$\chi_{0s}(r) \propto (1 - \exp[-(r - r_0)/\Lambda])$$

Protein-Water Interface

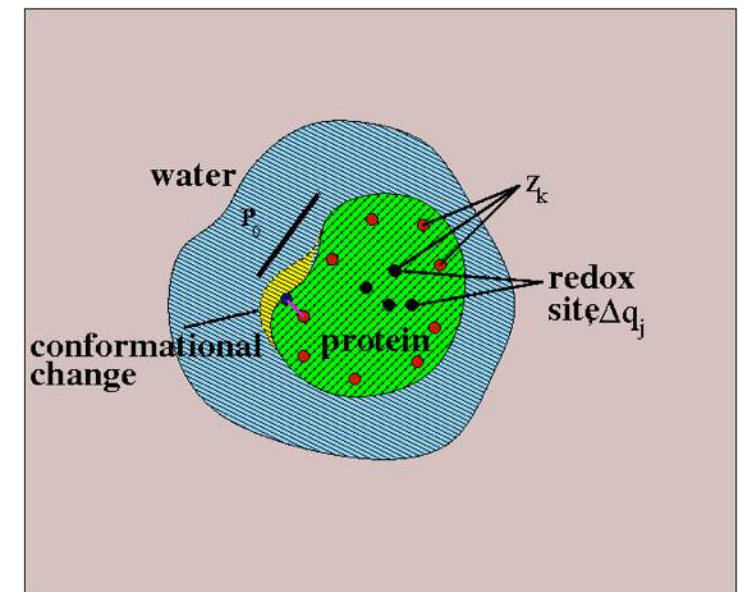
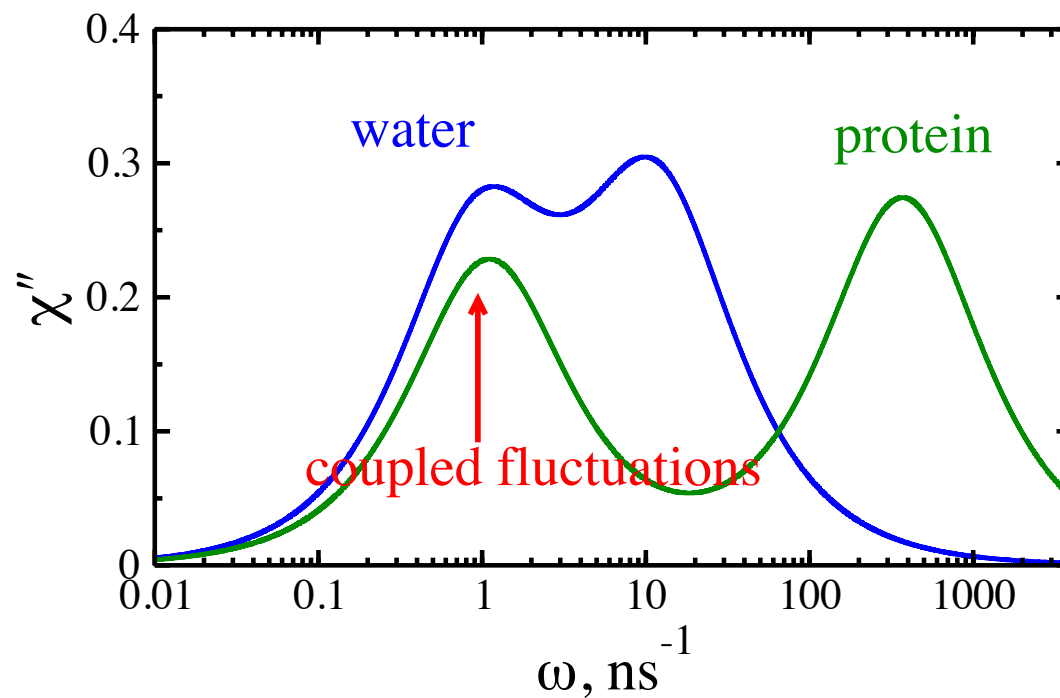
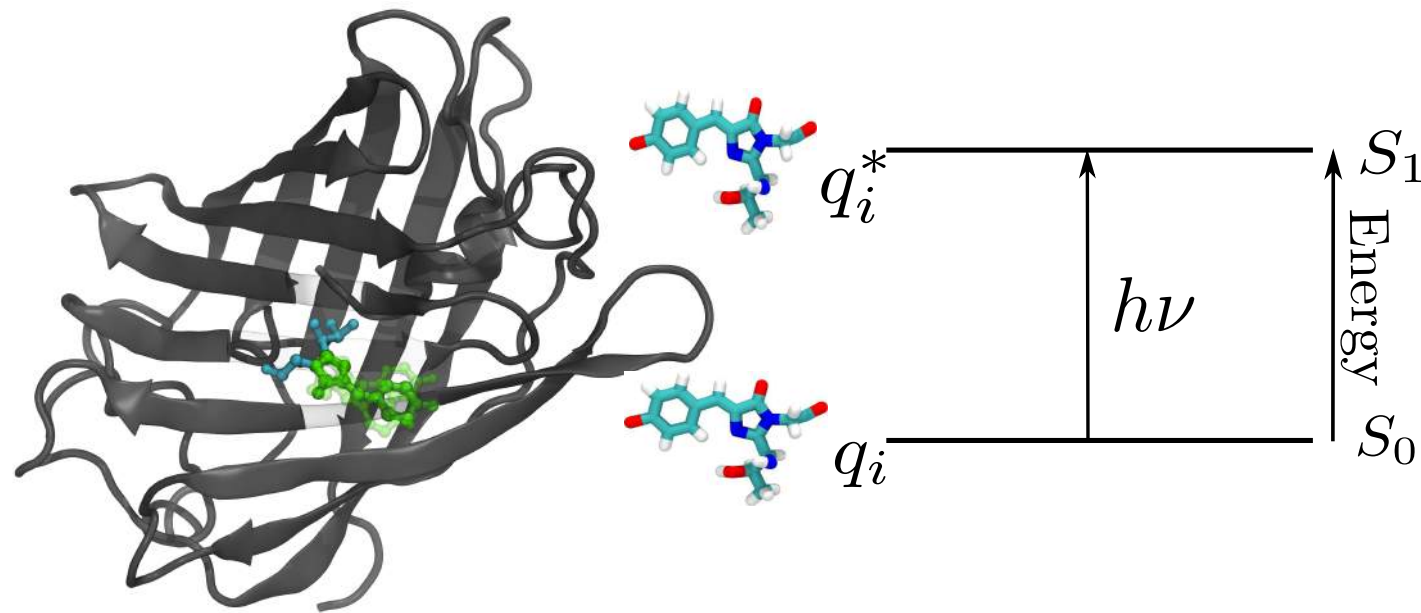


Water structure is locally broken: surface polarization is determined by the residue

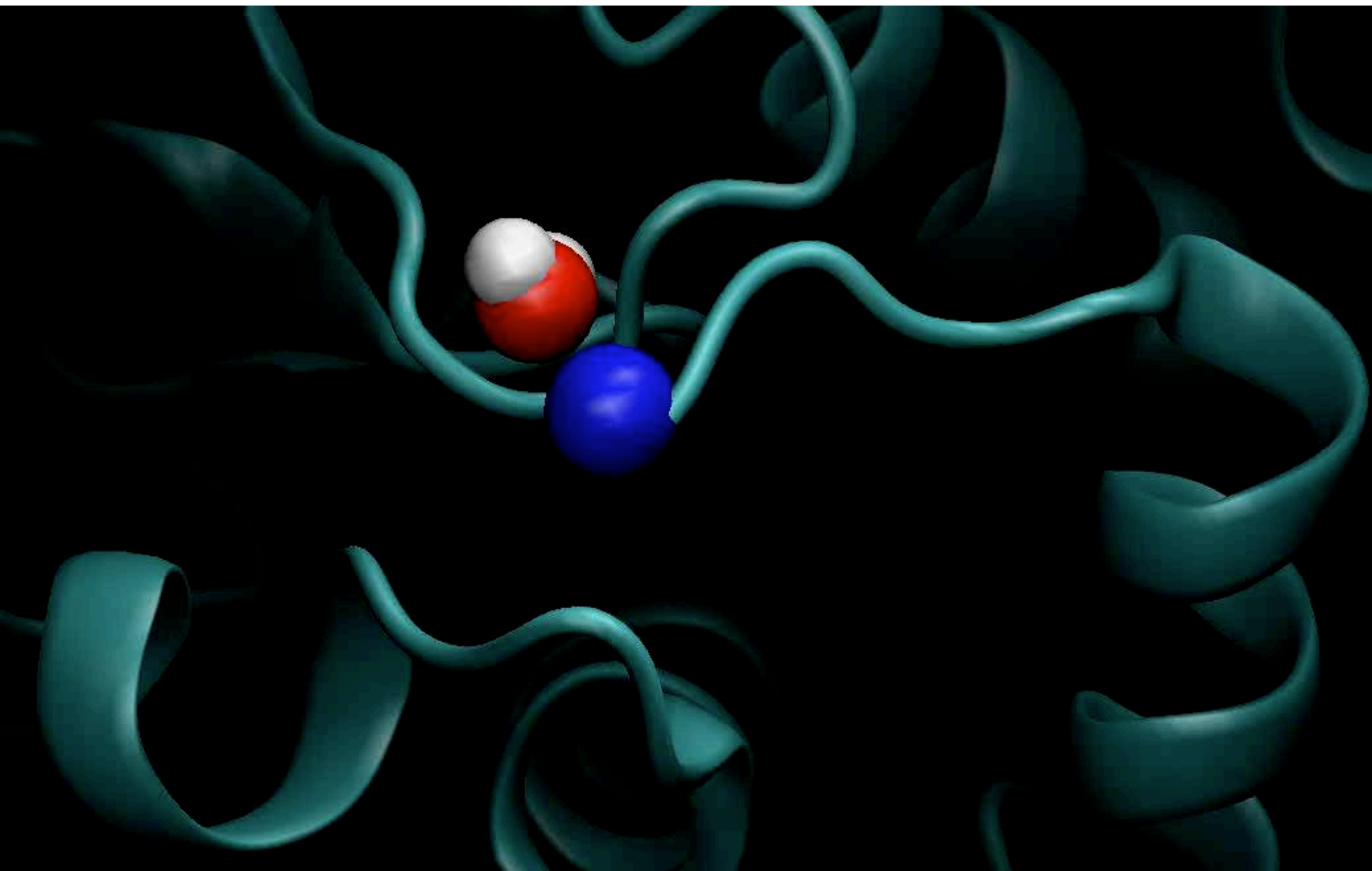
Frustration of surface polarized domains: long propagation into the bulk

Heterogeneous dynamics of the interfacial polarization

Coupled protein-water fluctuations (GFP)



Low-frequency motions of the protein move both the ionized surface residues and the water shells polarized by them.

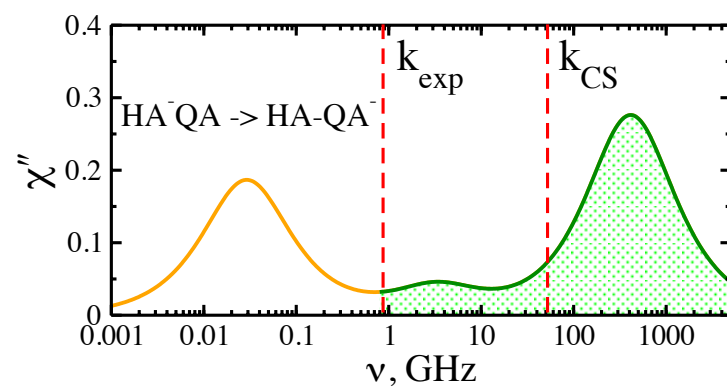


Proteins:

Large-scale electrostatic fluctuations (lots of surface charges)

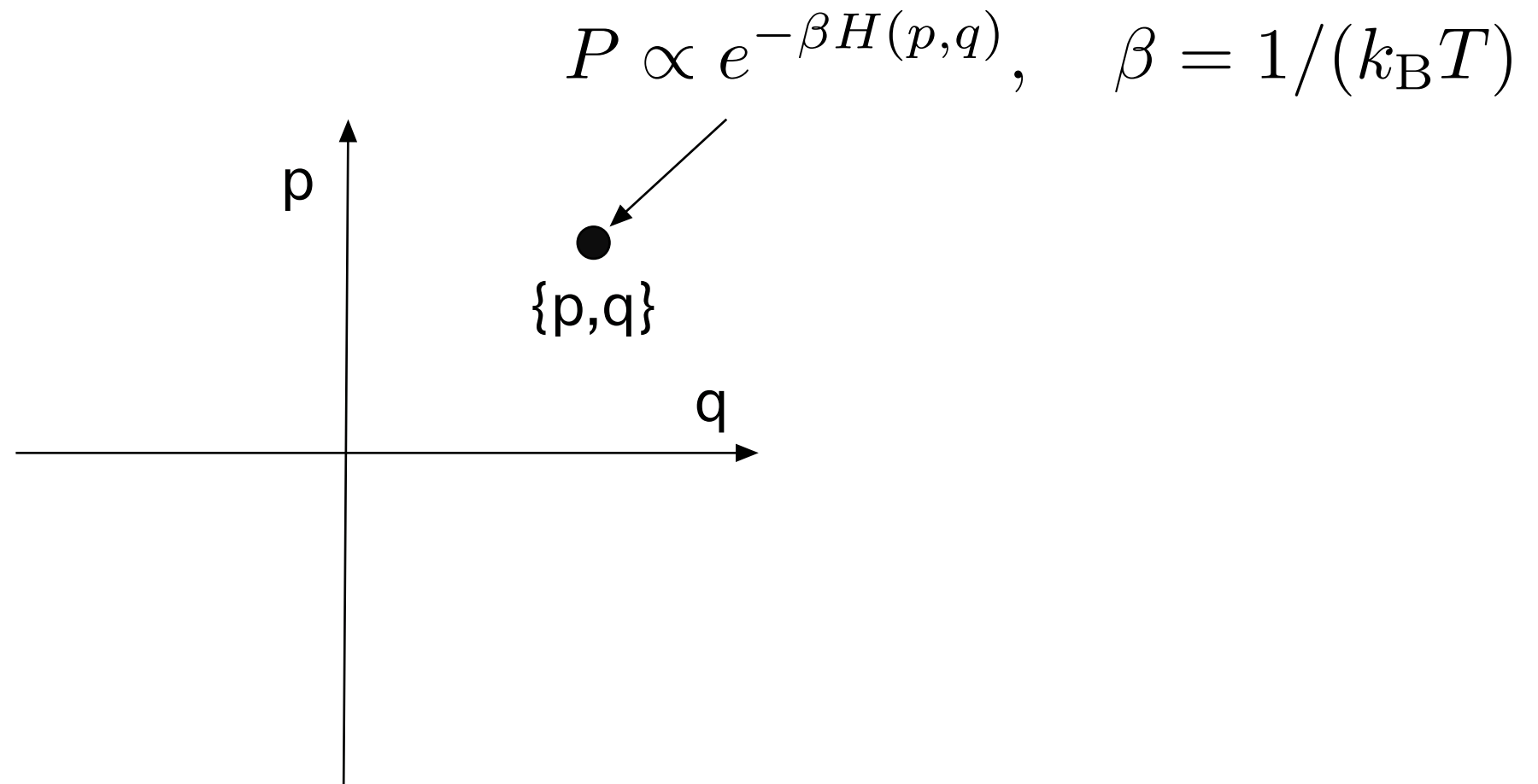
Slow modes

What if the rate of the reactions is faster than the fluctuations?



?

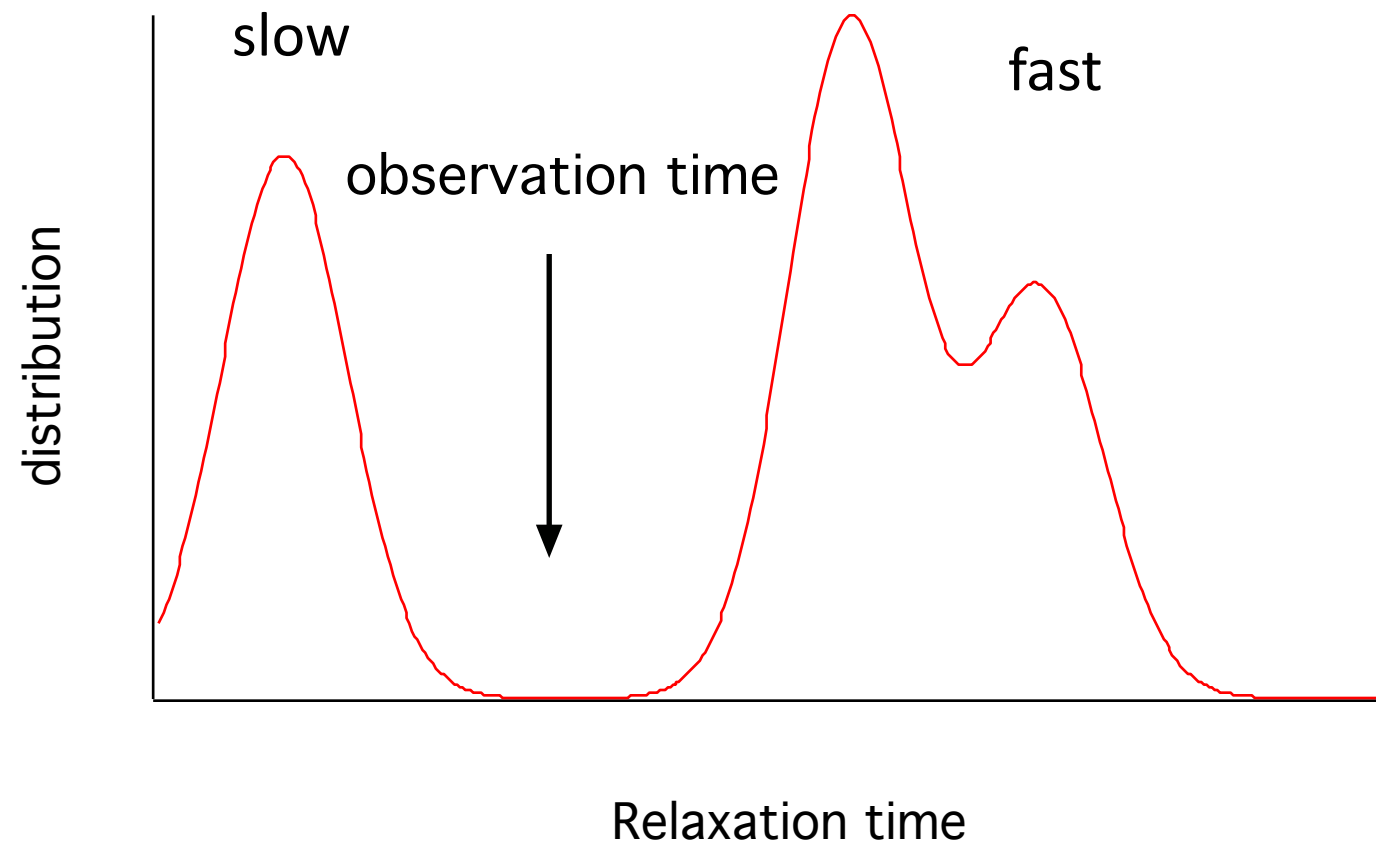
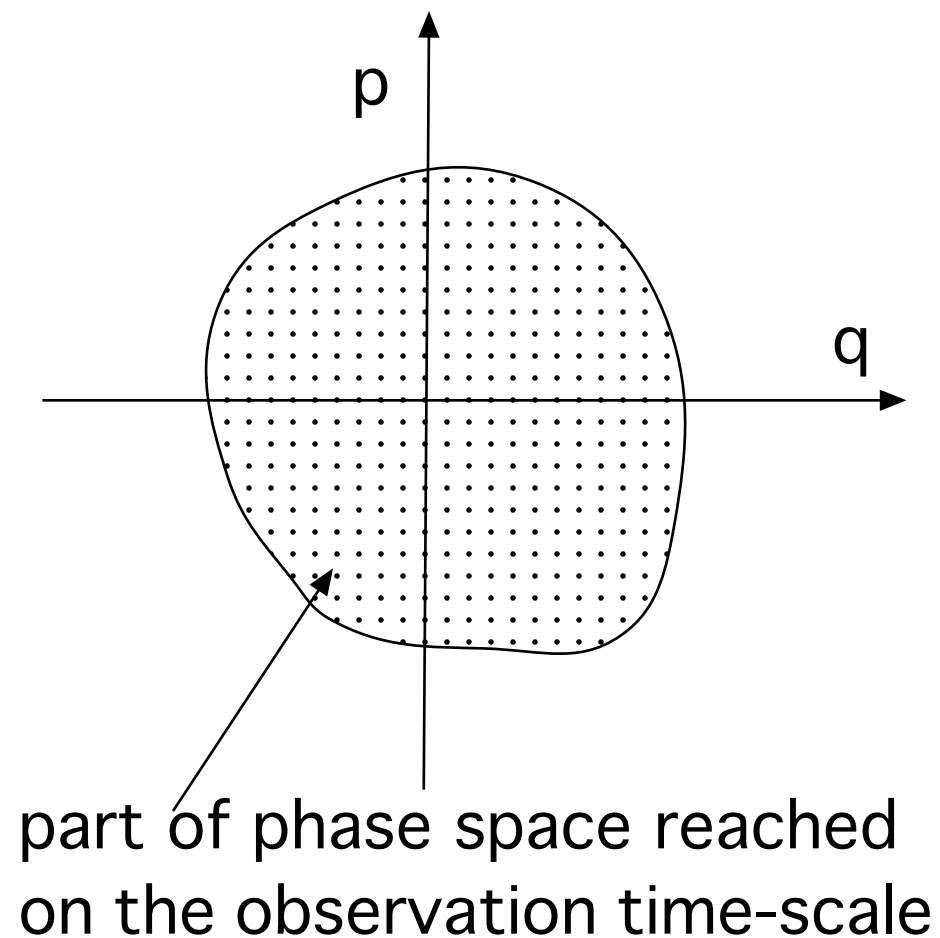
Ergodicity



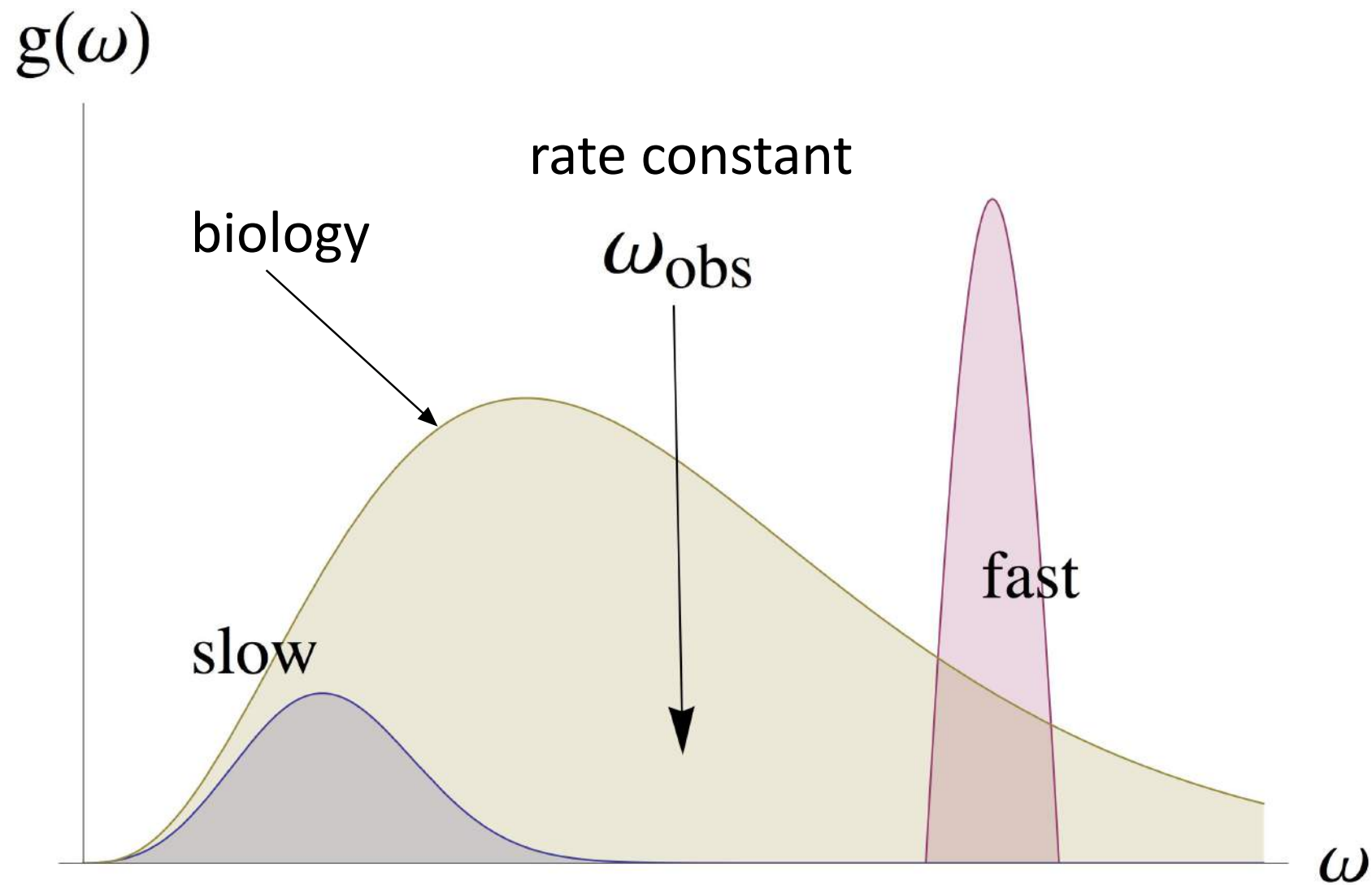
Mathematical abstraction: $\tau_{\text{obs}} \rightarrow \infty$

Canonical ensemble: “All the ‘fast’ things have happened and all the ‘slow’ things have not”

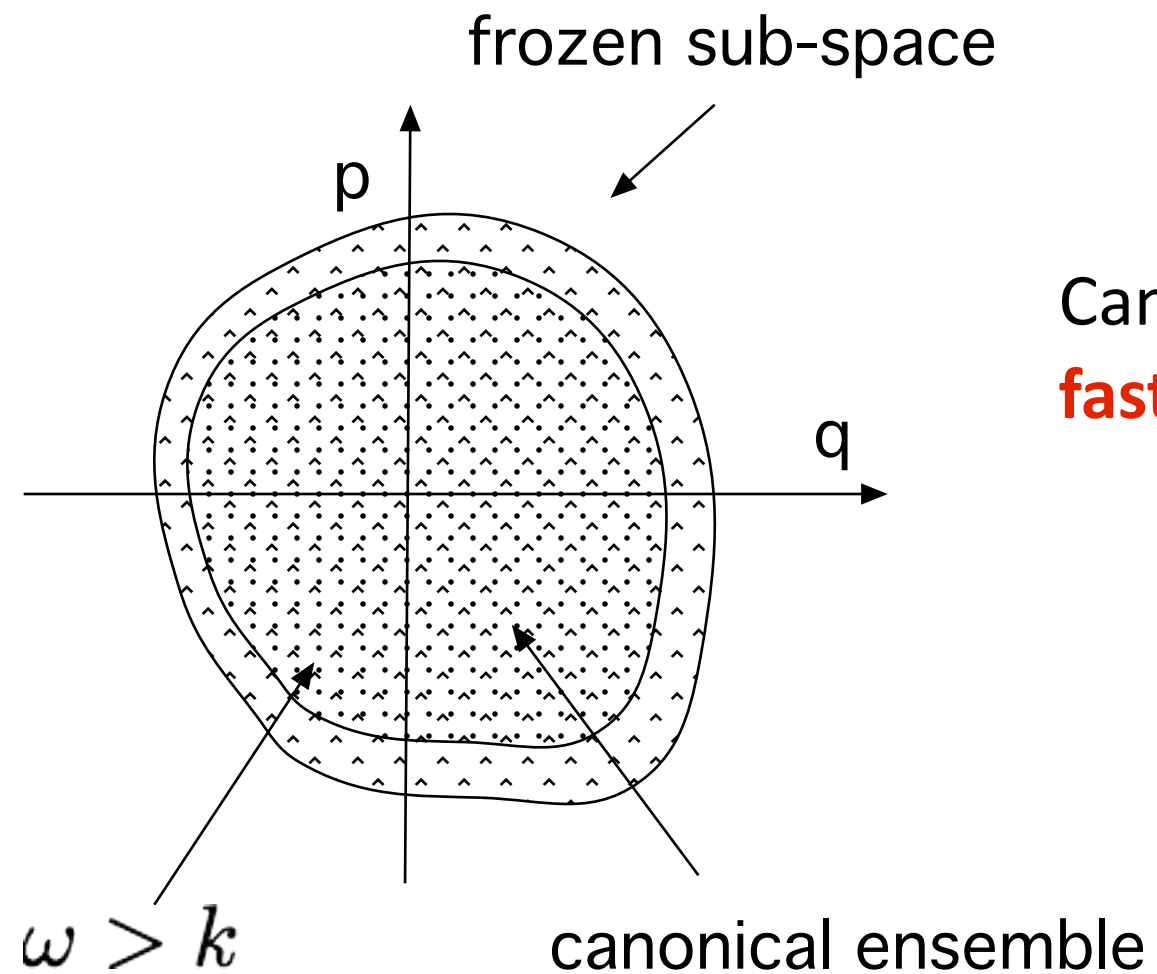
- R. Feynman



Biology: continuous ergodicity breaking



Dynamically **restricted** ensemble

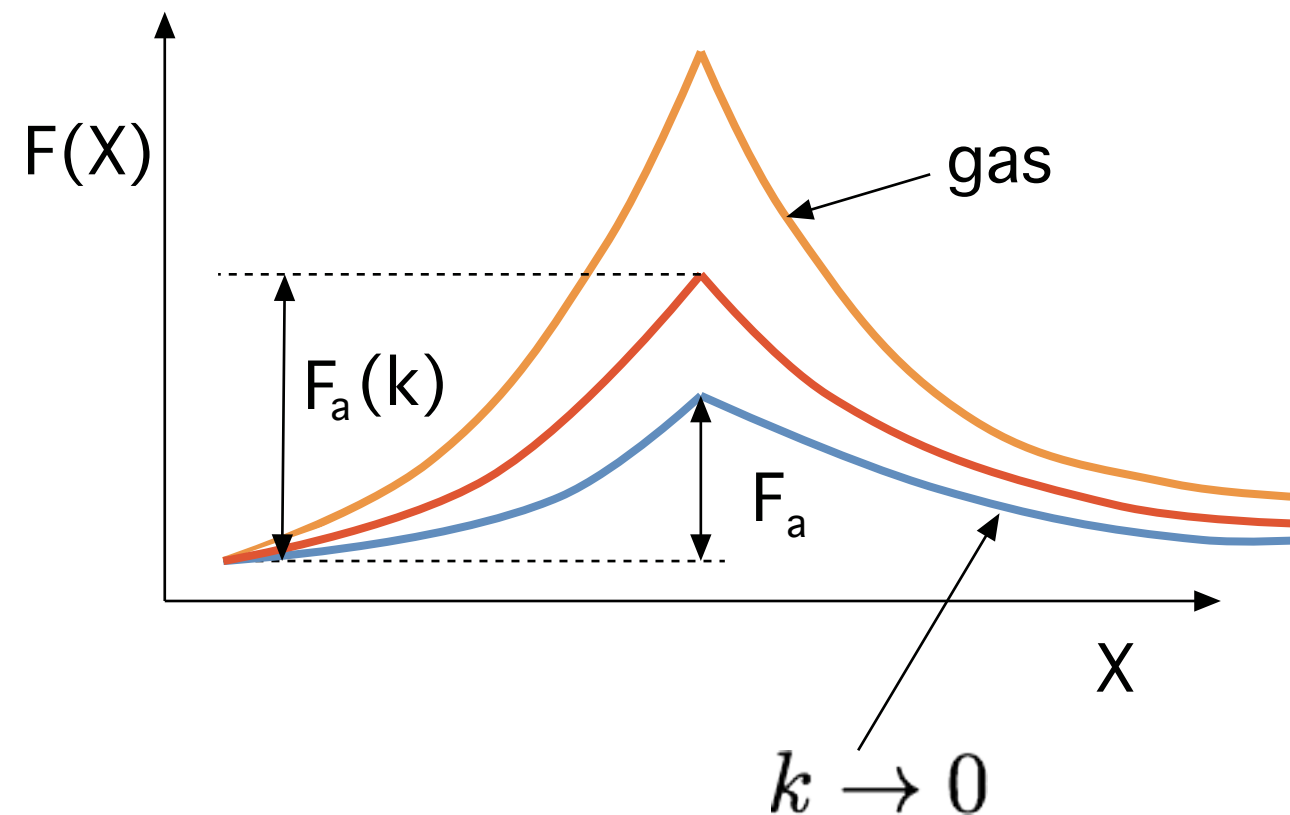


Canonical (Gibbs) average over motions
faster than the rate

$$k < \omega$$

Nonergodic kinetics

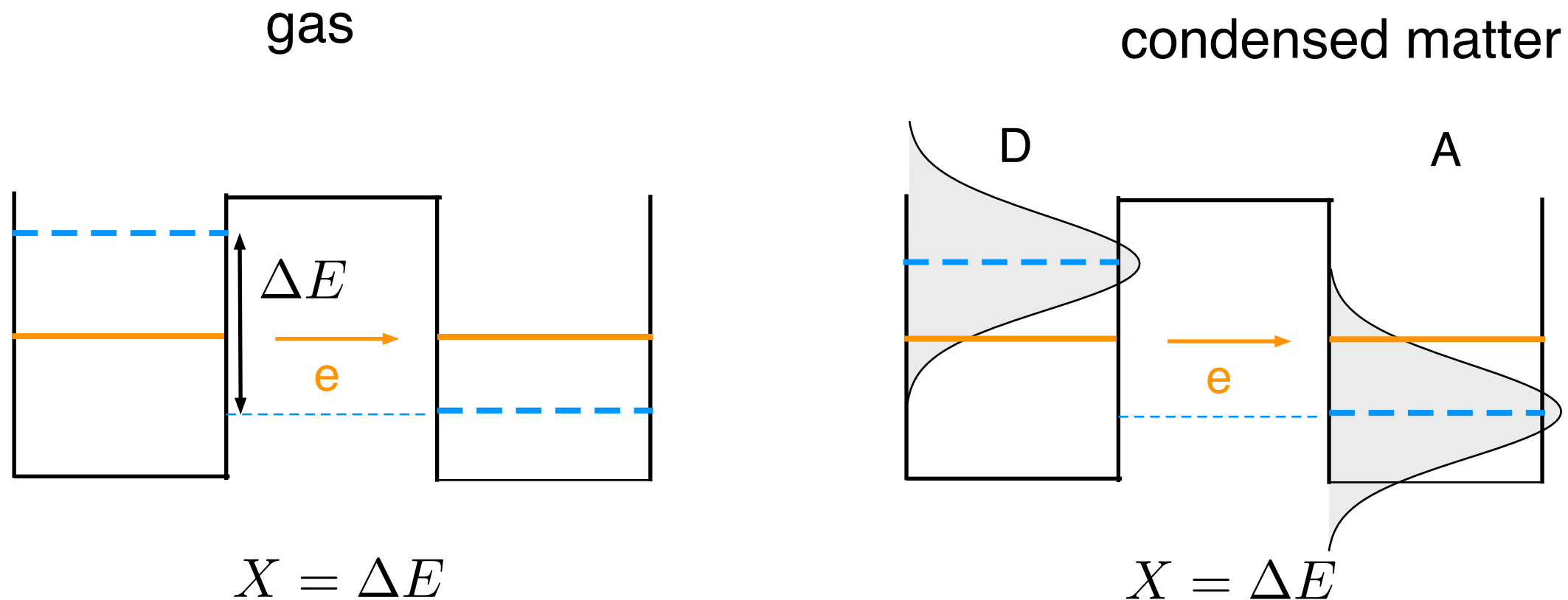
Activation barrier depends on the rate



$$\mathbf{k} \propto \exp[-\beta F_a(\mathbf{k})]$$

self-consistent solution for \mathbf{k}

Electron tunneling (hopping conductivity)



$$F(X) = -k_{\text{B}}T \ln P(X) = \frac{(X - X_0)^2}{4\lambda}$$

Reorganization energy (extent of medium deformation)

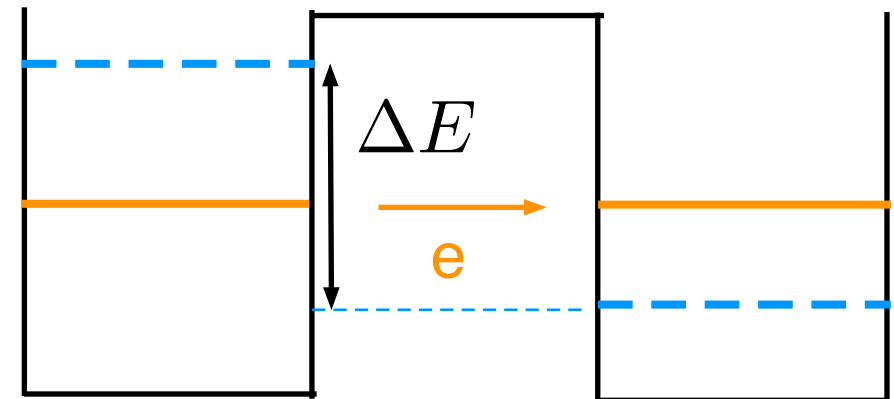
Dynamics vs thermodynamics

Stokes-shift time correlation function:

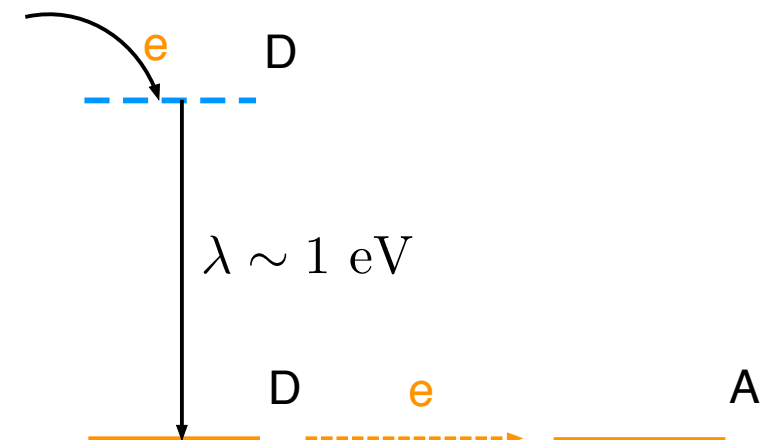
$$C_i(t) = \langle \delta X(t) \delta X(0) \rangle_i$$

$$\chi_i''(\omega) = (\beta\omega/2)C_i(\omega)$$

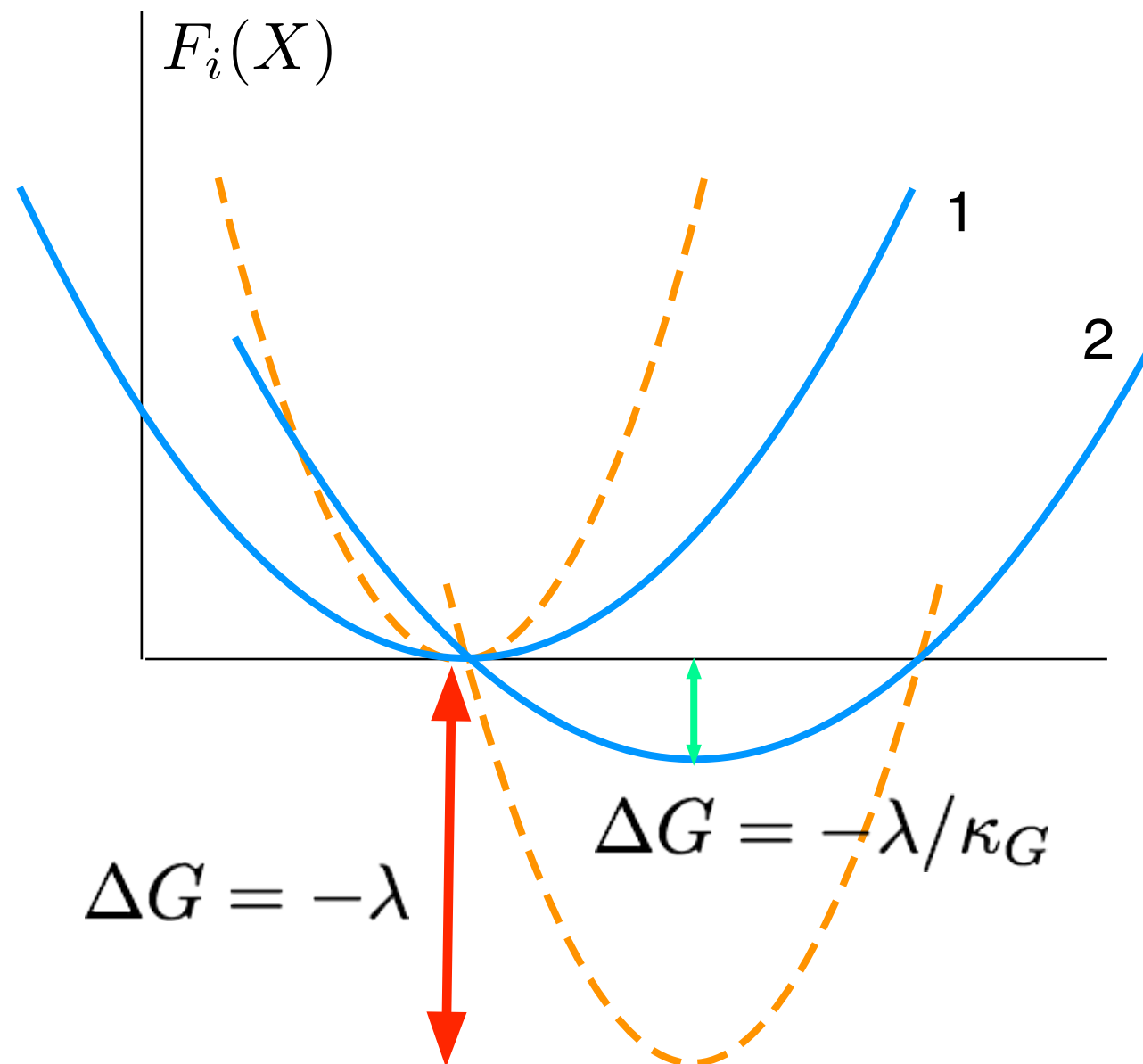
$$\lambda_s(k) \propto \int_{k=\omega_{\text{obs}}}^{\infty} \chi''(\omega) (d\omega/\omega)$$



$$X = \Delta E$$

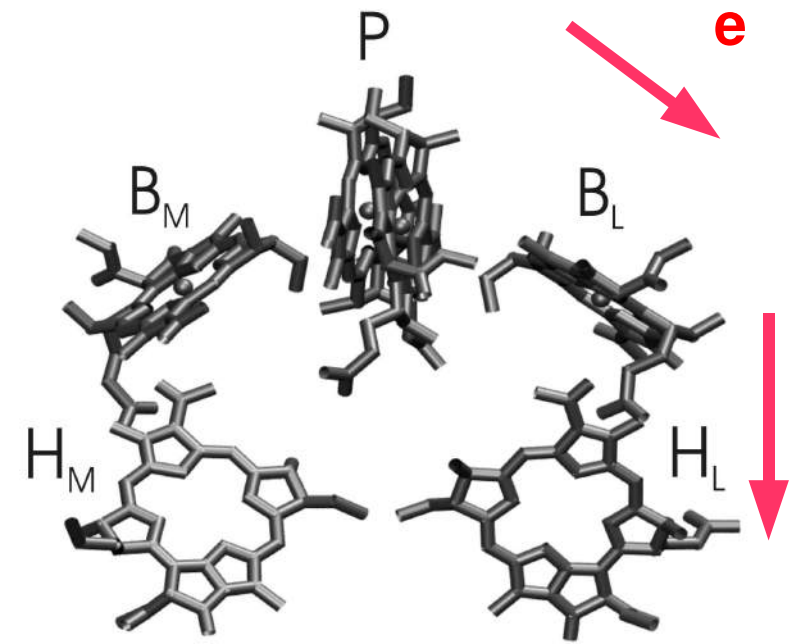
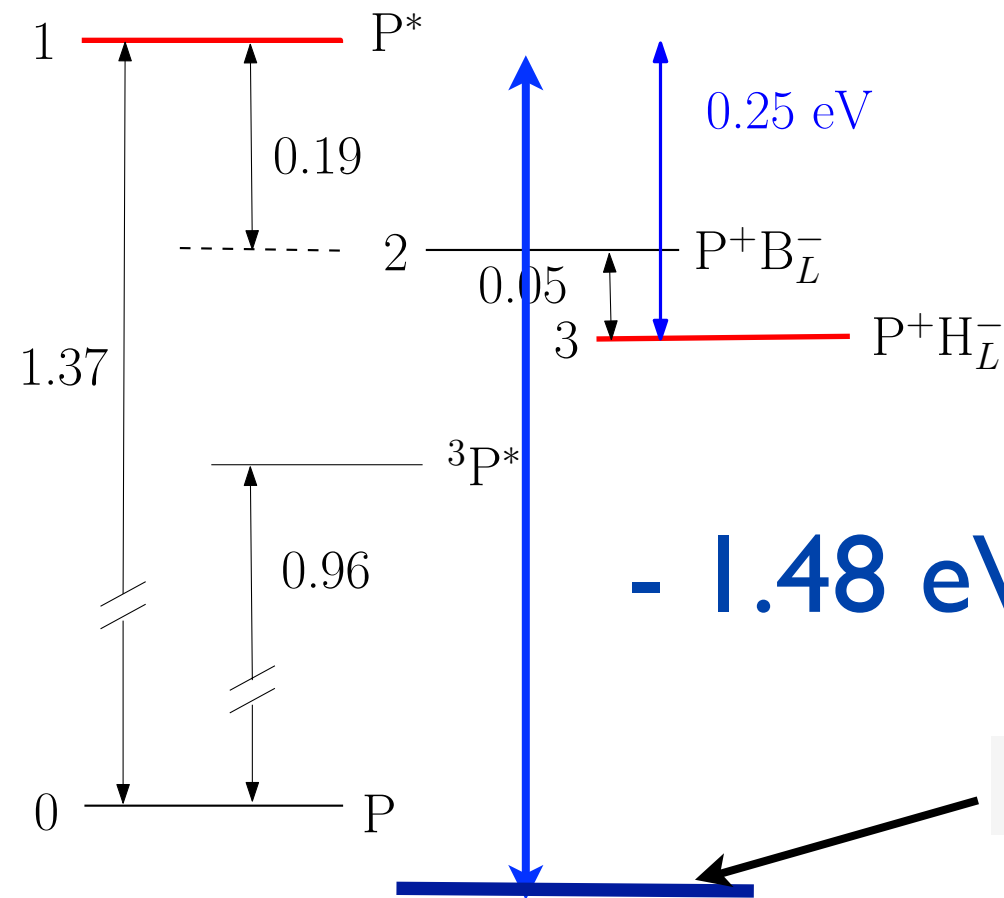


Reaction free energy & activation barrier



$\kappa_G \sim 3 - 20$ from MD of redox proteins

Bacterial photosynthesis: Energetics

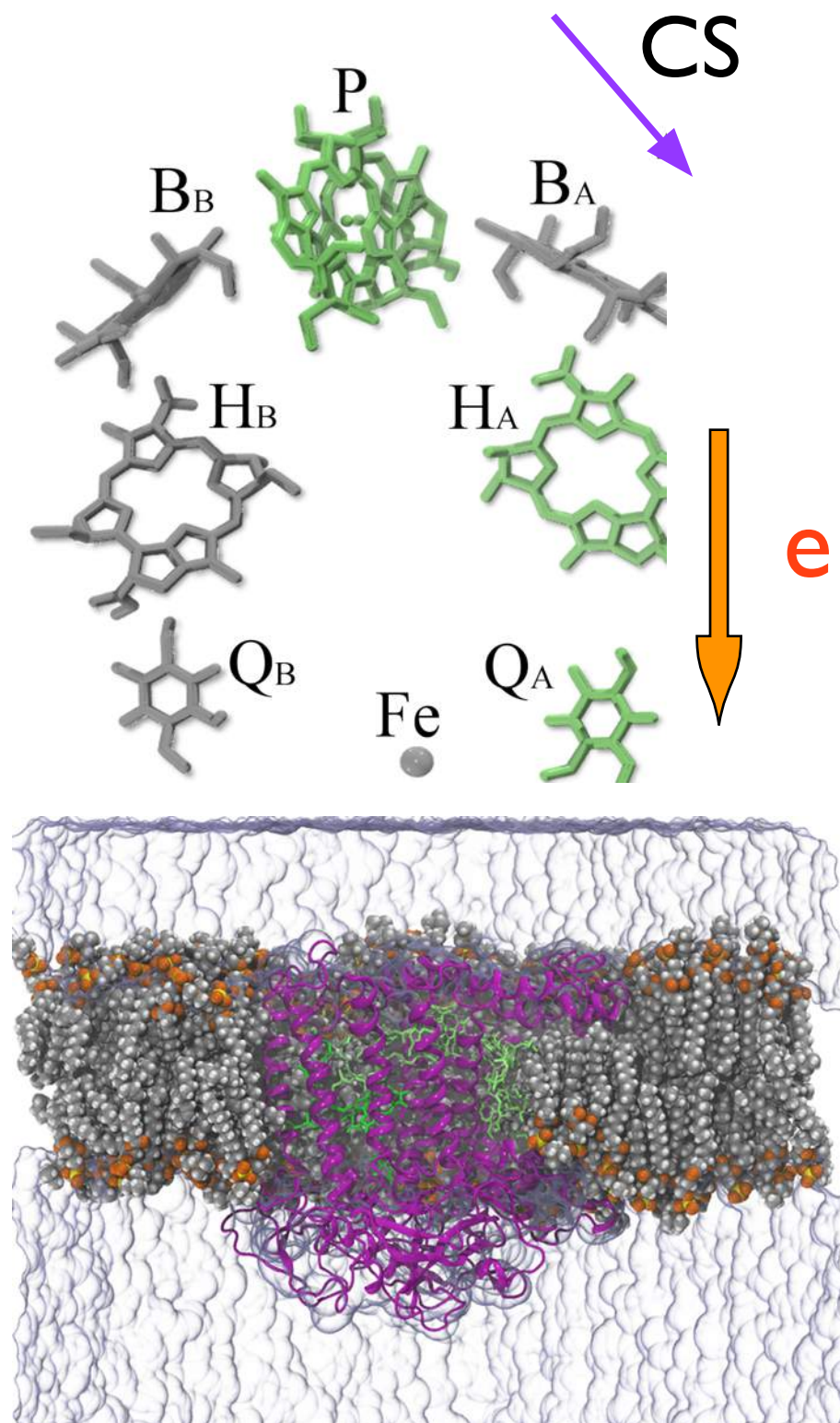
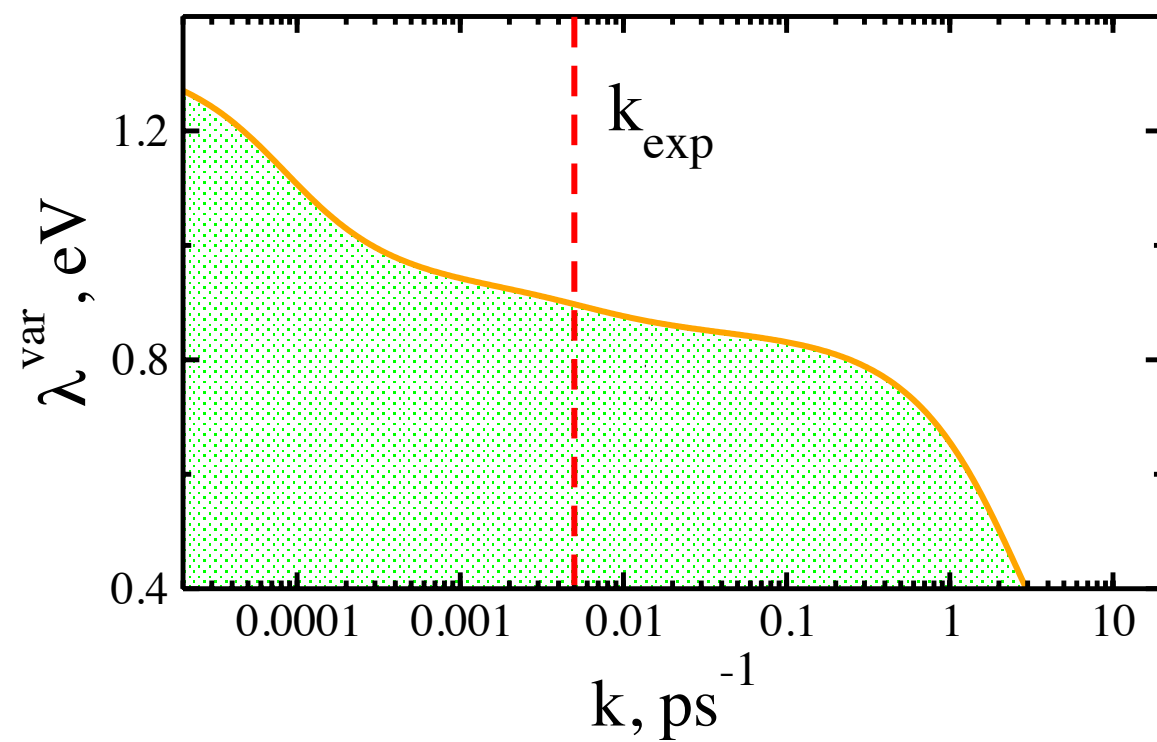
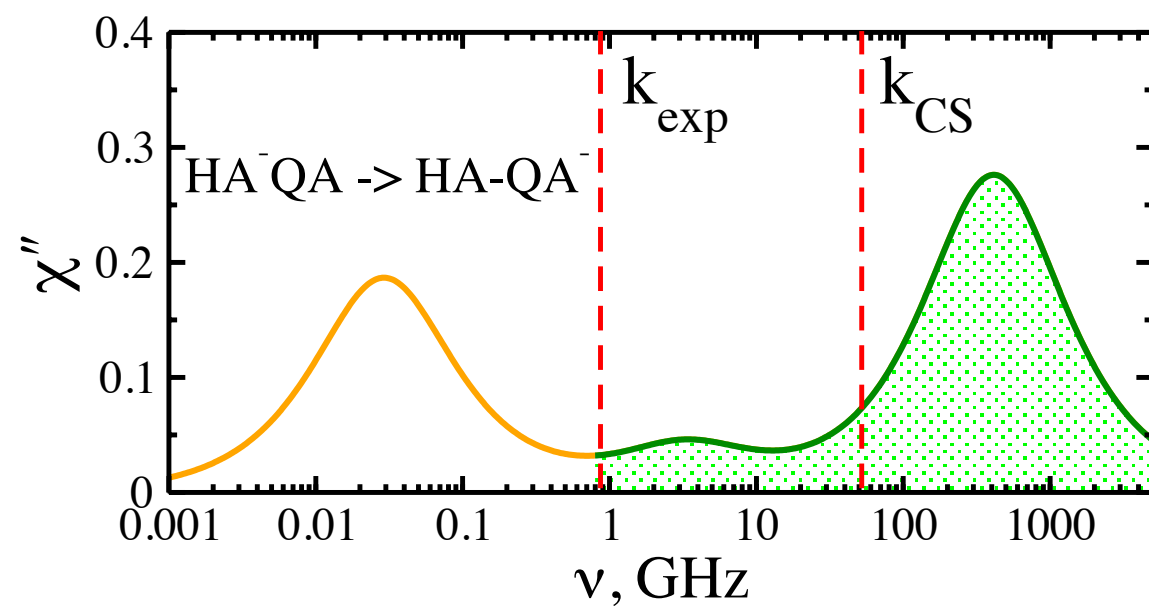


traditional theory

$$\lambda_{12} \simeq 0.77 \text{ eV}, \lambda_{23} = 0.71 \text{ eV}$$

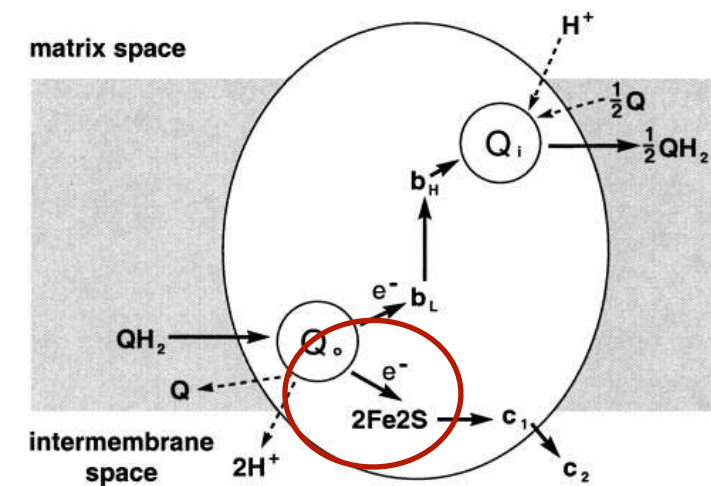
$$-\Delta G_{13} = \lambda_{12} + \lambda_{23} = 1.48 \text{ eV}$$

HA to QA

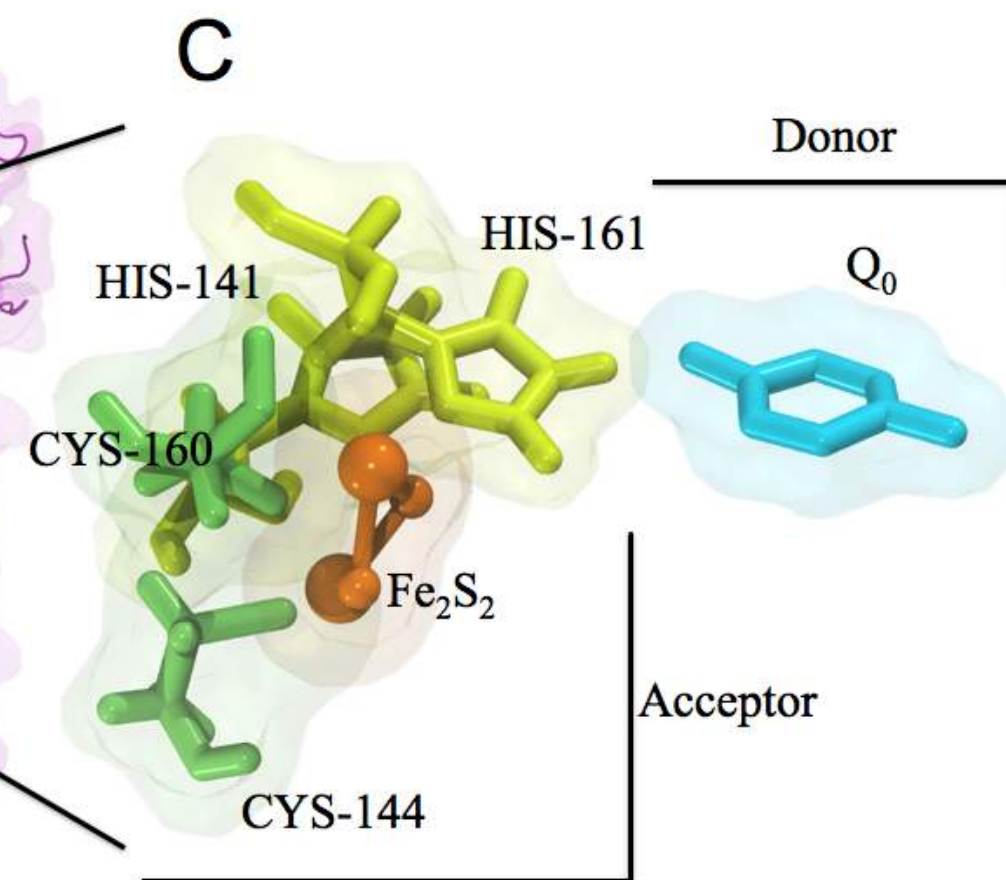
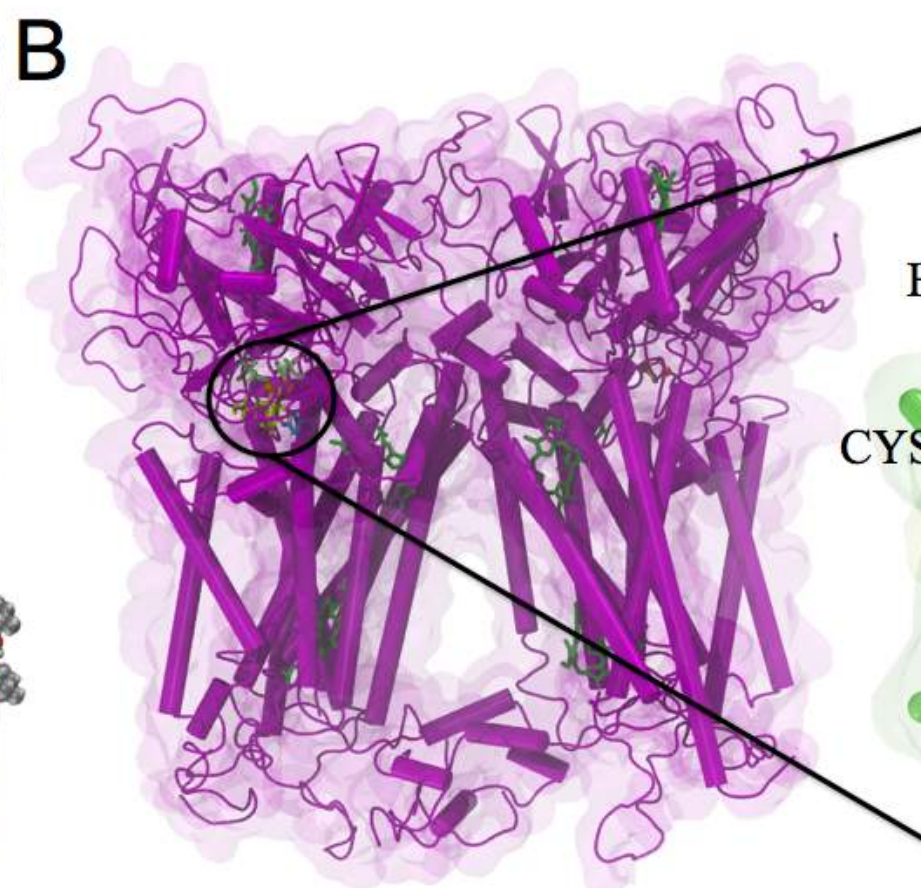
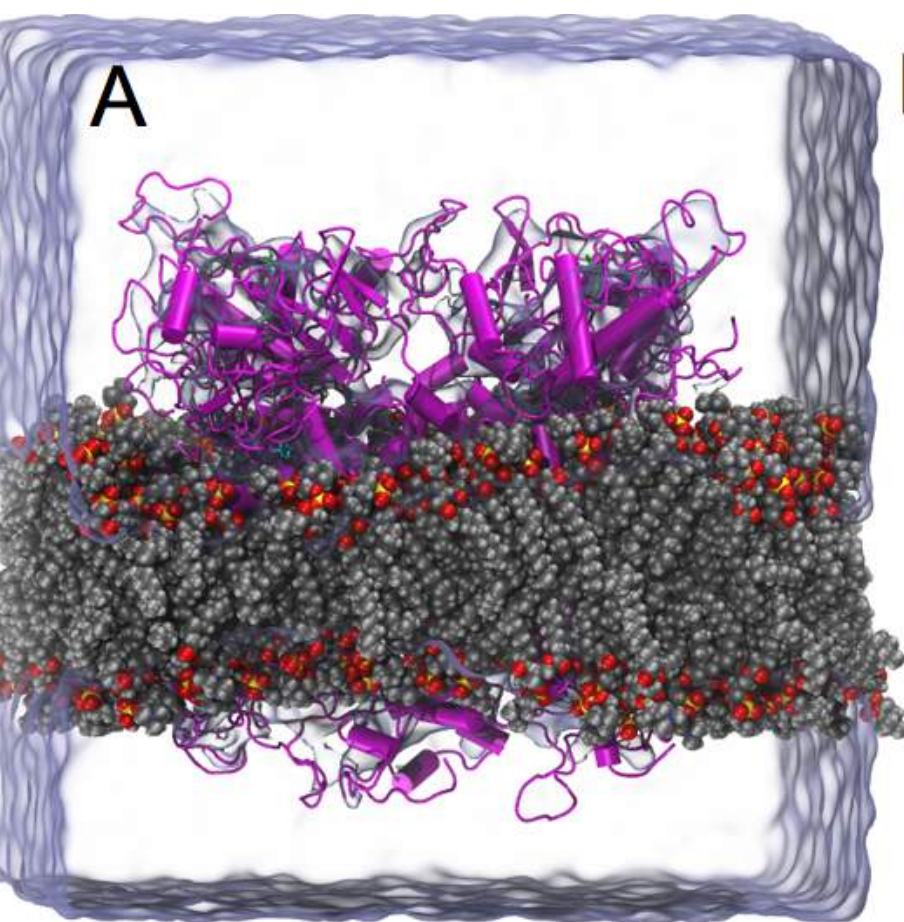


bc1 complex of respiratory chains

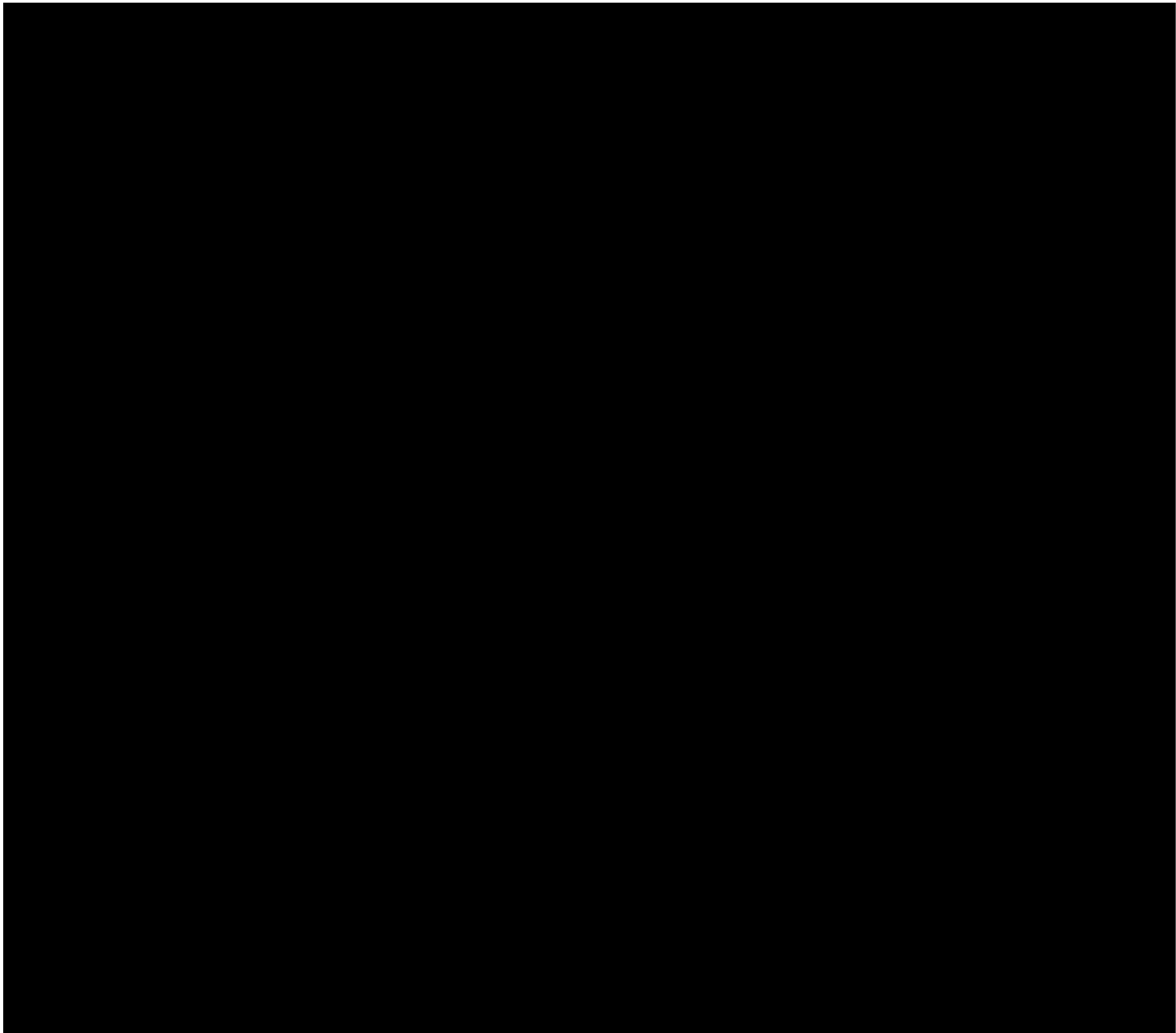
Reaction rate ~ 1 ms



Izrailev et al., Biophys. J. '99

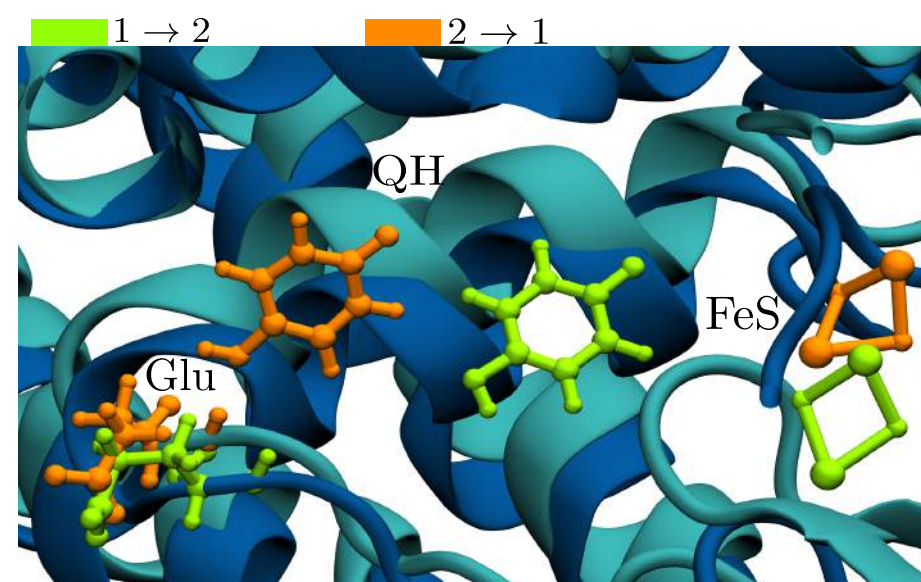
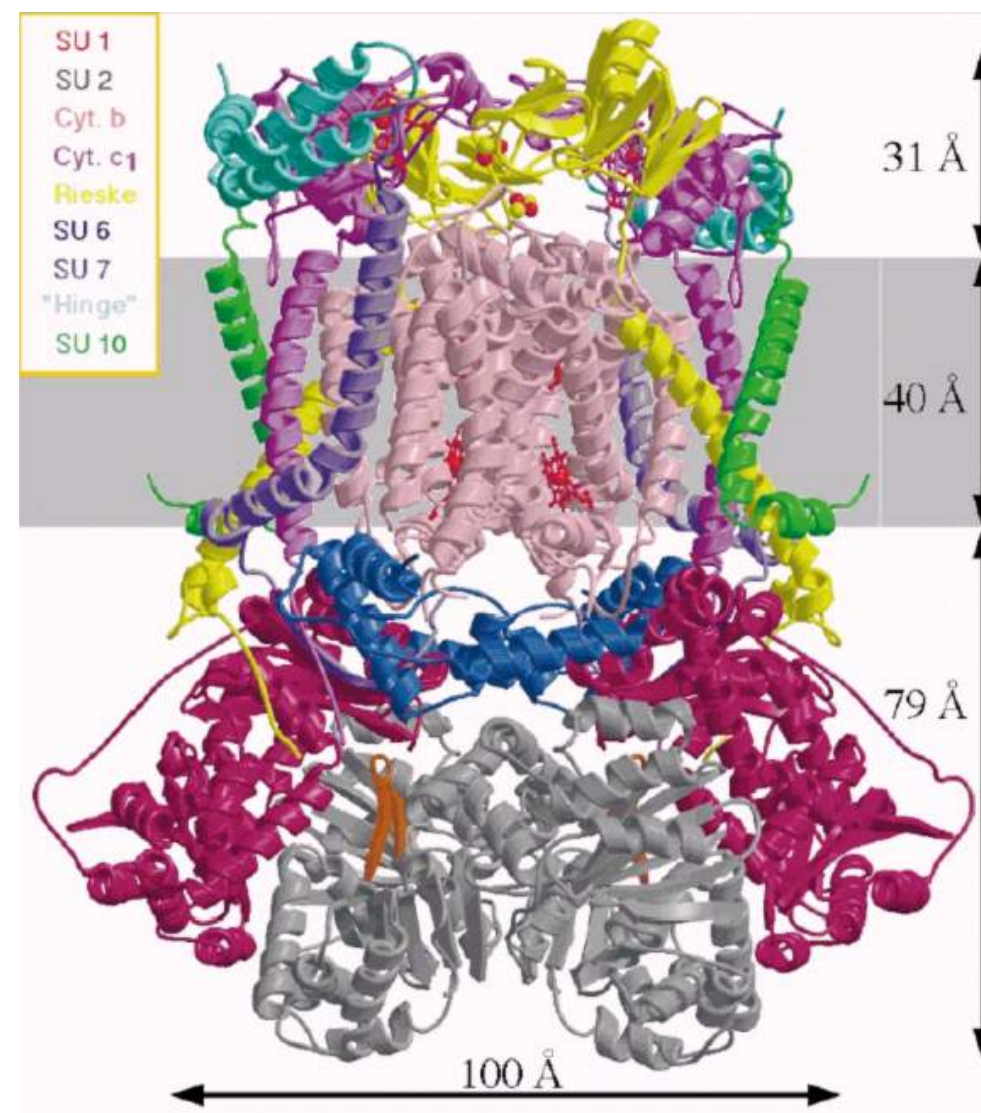
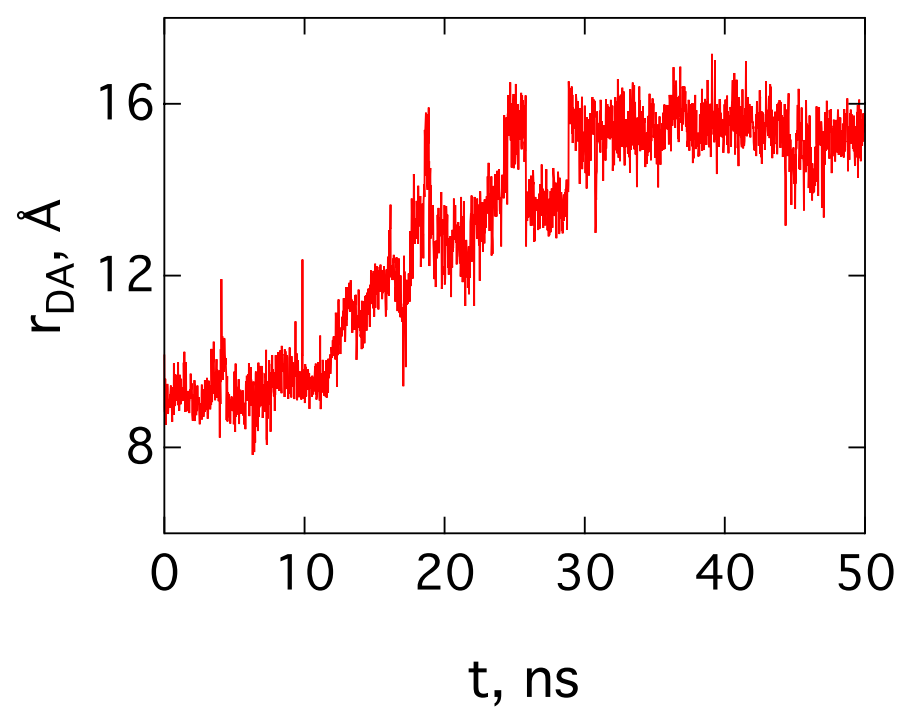
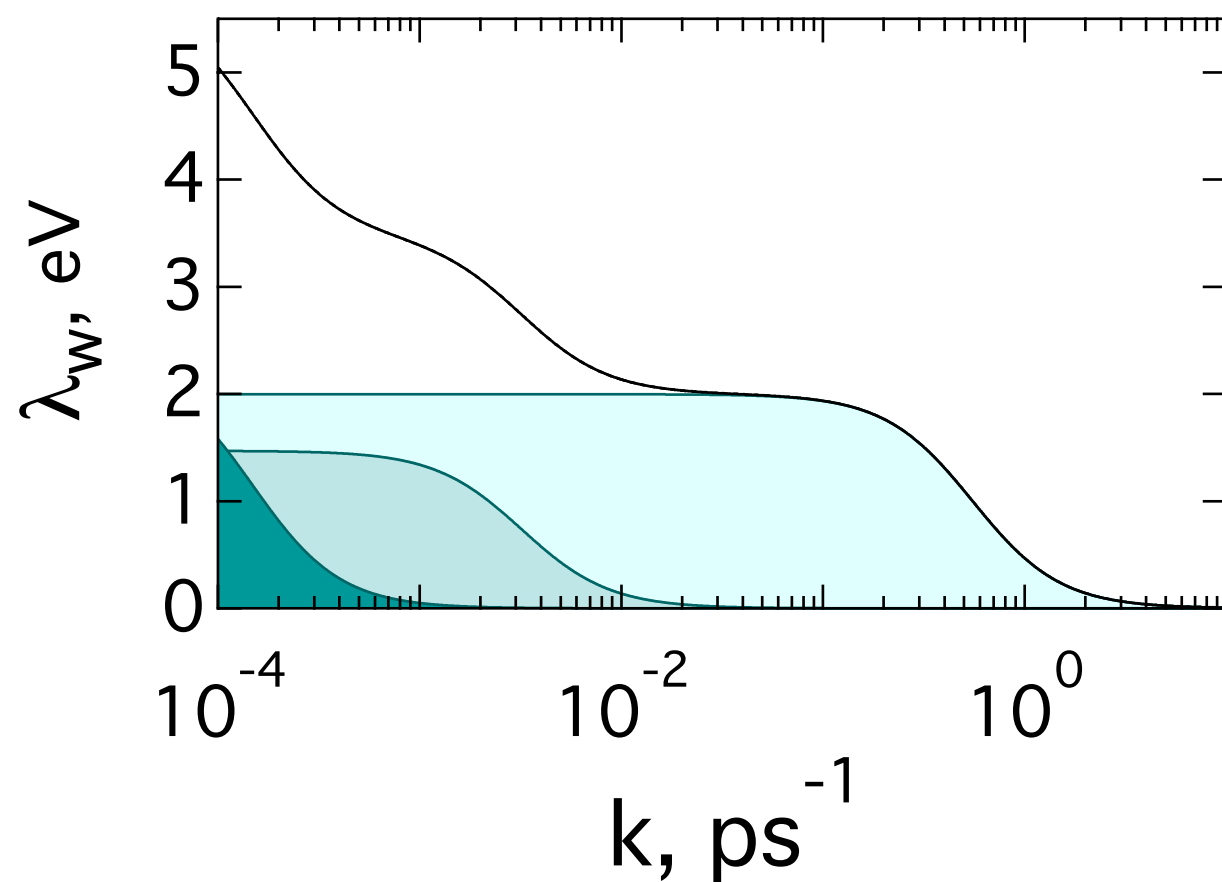


Martin, LeBard, DVM, JPCL 4 (2014) 3602.



bc1 complex: $\lambda(k)$

200 ns simulation

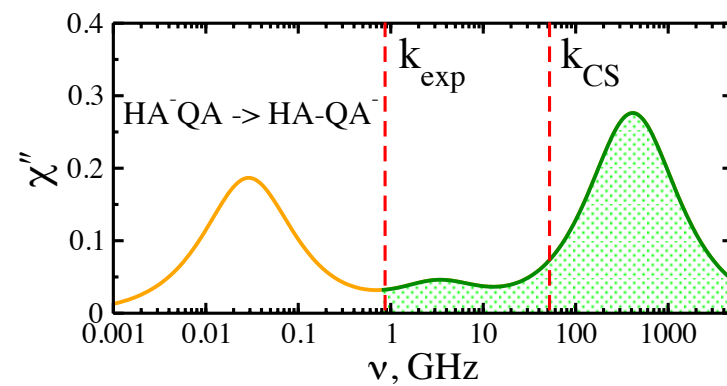


Why are enzymes big?

We are all familiar with the systems like software (or government legislation) ... can grow rapidly in size over a number of years. Enzyme evolution is a great deal slower - but it has been going for million of years...

"From Enzyme Models to Model Enzymes", Kirby & Hollfelder

Maybe because slow reactions require slower elastic deformations of the interface



...the advancement of science depends on the discovery and development of exact ideas ...
to warrant the deductions we may draw by the application of mathematical reasoning.

- J. C. Maxwell

Biological interfaces are “different”

Nonergodicity introduces timescales to where previously
only (free) energy has ruled.



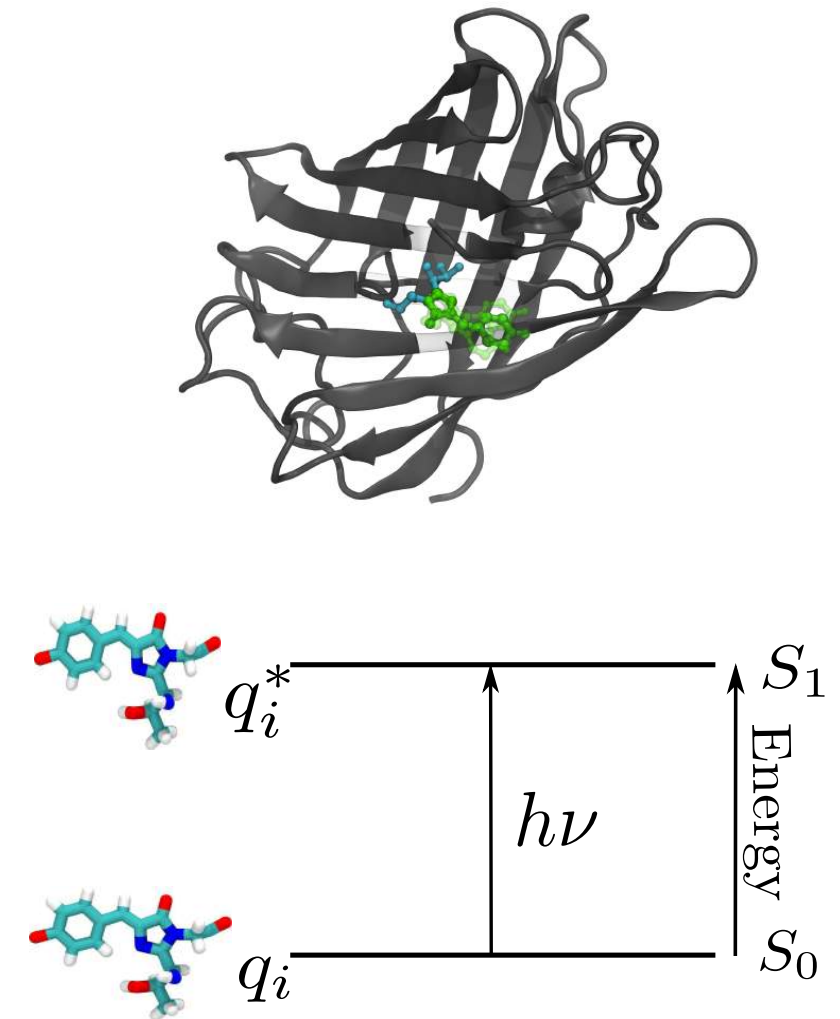
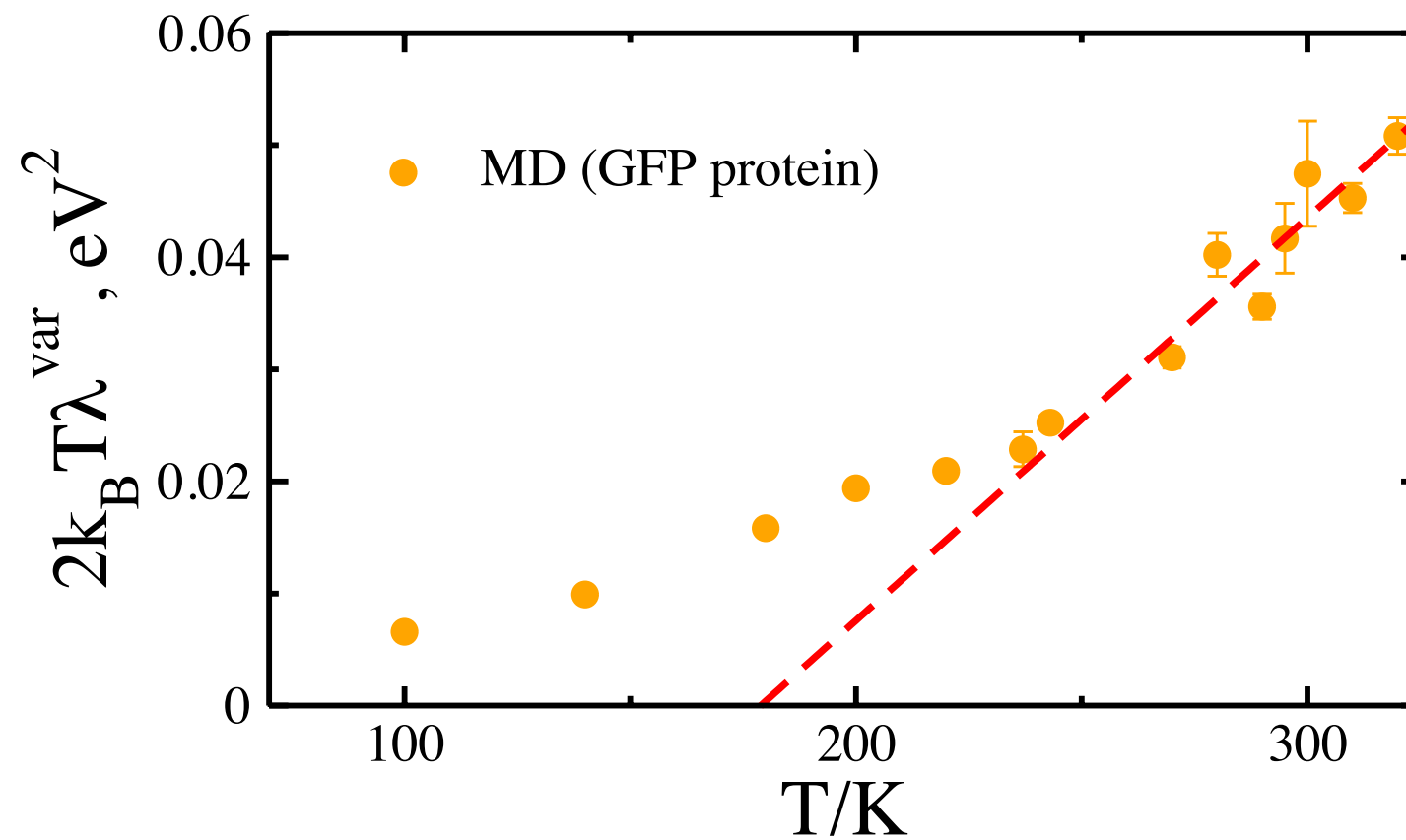
Prof David LeBard



Dr Daniel Martin

\$\$ NSF

Lambda: dynamical transition

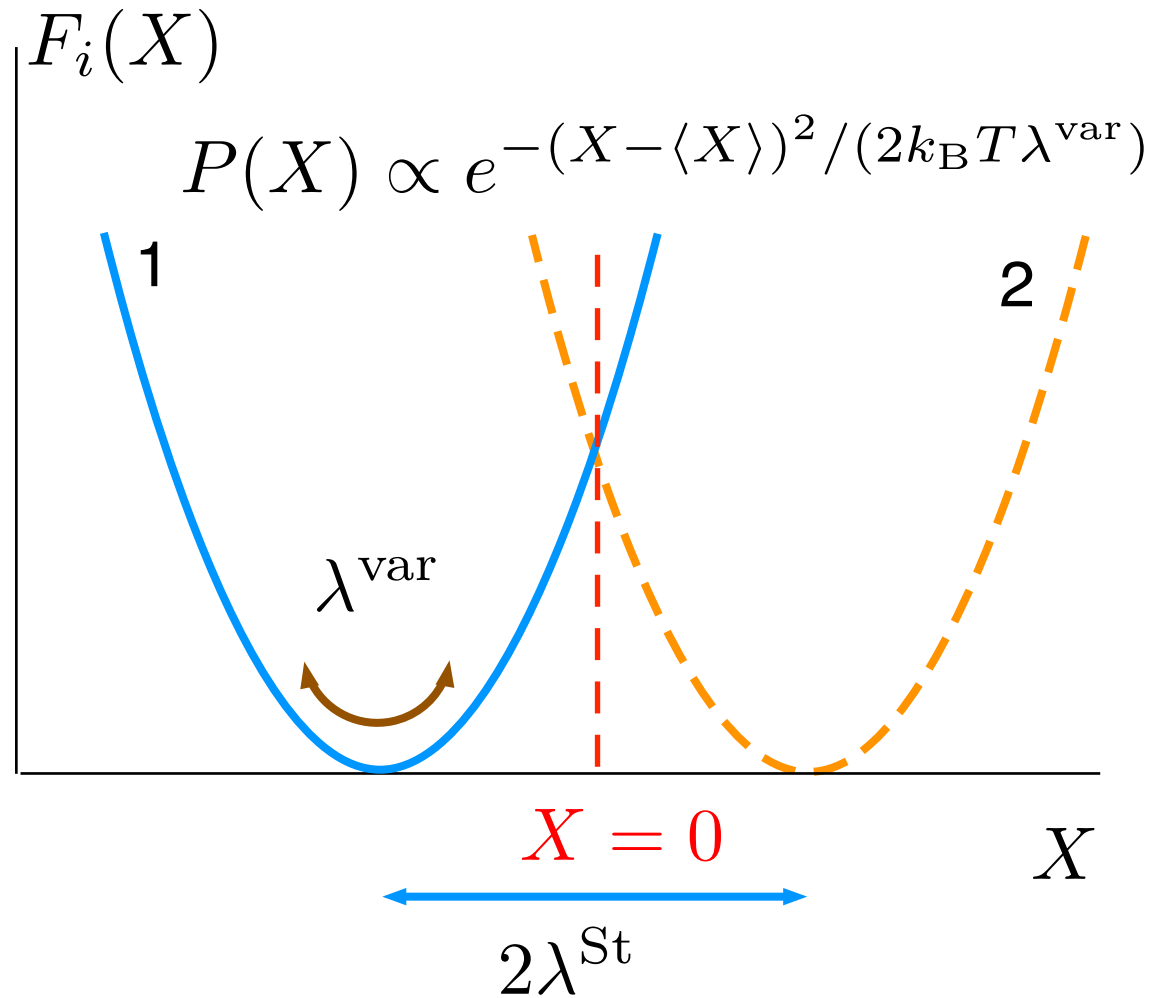


Standard picture: $\sigma(T)^2 = 2\lambda^{\text{var}} k_B T \propto T$

Proteins: $2k_B T \lambda^{\text{var}} \propto T - T_0$

JPCB 116 (2012) 10294

Glassy kinetics



$$k_{\text{ET}} \propto P(0)$$

$$2k_B T \lambda^{\text{var}} \propto T - T_0$$

High-temperature rate constant:

$$k_{\text{ET}} \propto \exp \left(-\frac{A}{T - T_0} \right)$$

Fogel-Fulcher-Tammann high-temperature kinetics

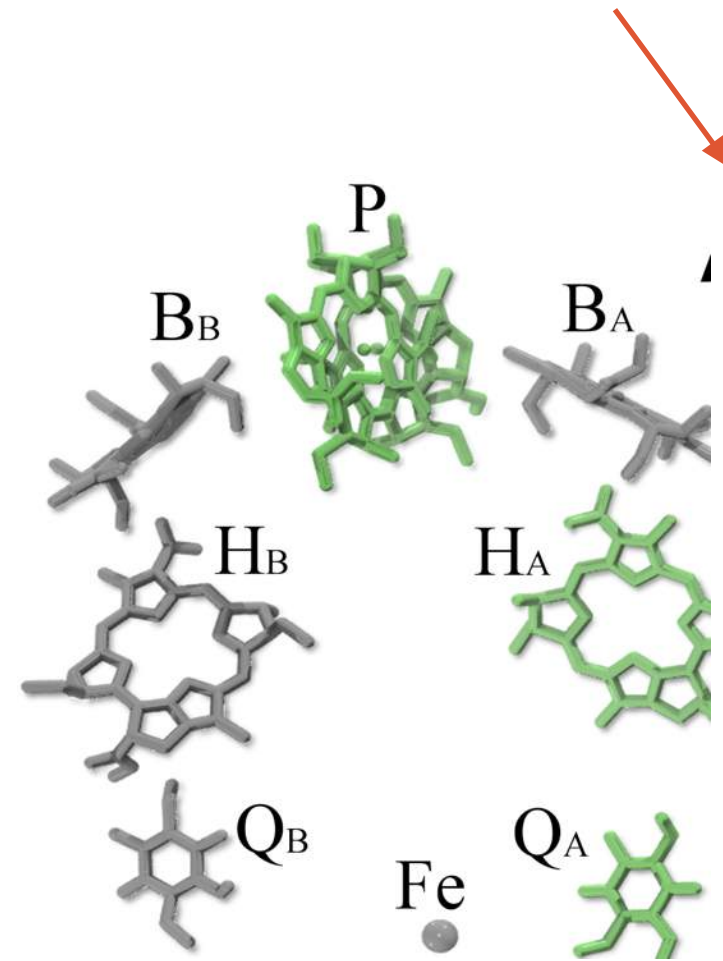
Bacterial charge separation (3 ps reaction time!)

$$\lambda(k) \propto \int_k^\infty \chi''(\omega)(d\omega/\omega)$$

$$\underbrace{\mathbf{k} \propto \exp[-\beta F_a(\mathbf{k})]}_{\text{self-consistent solution for } \mathbf{k}}$$

$$\lambda(k = 0.3 \text{ ps}^{-1}) = 0.35 \text{ eV}$$

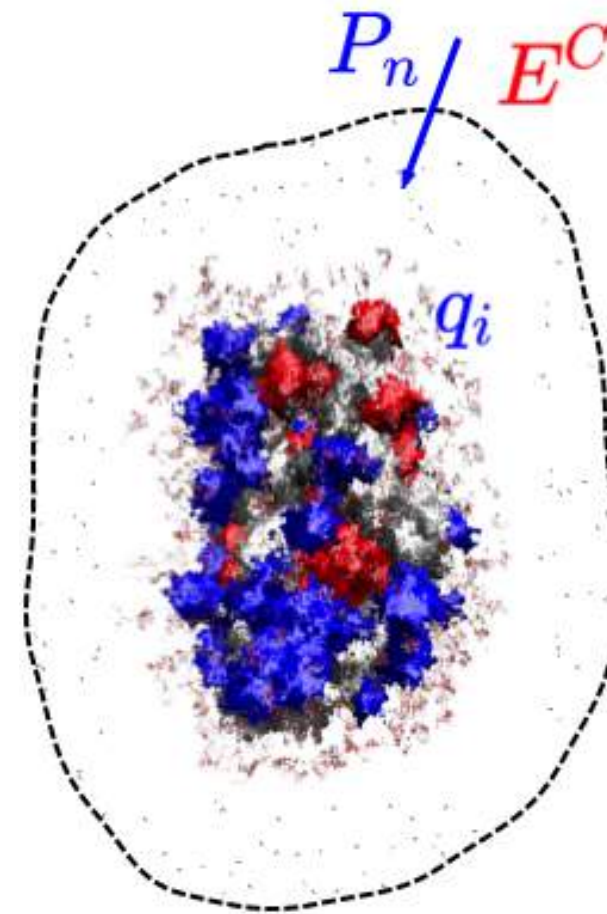
$$\lambda(k \rightarrow 0) = 2.4 \text{ eV}$$



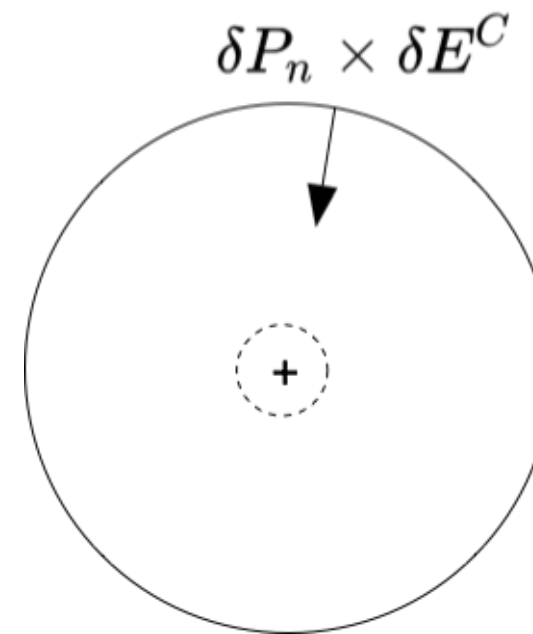
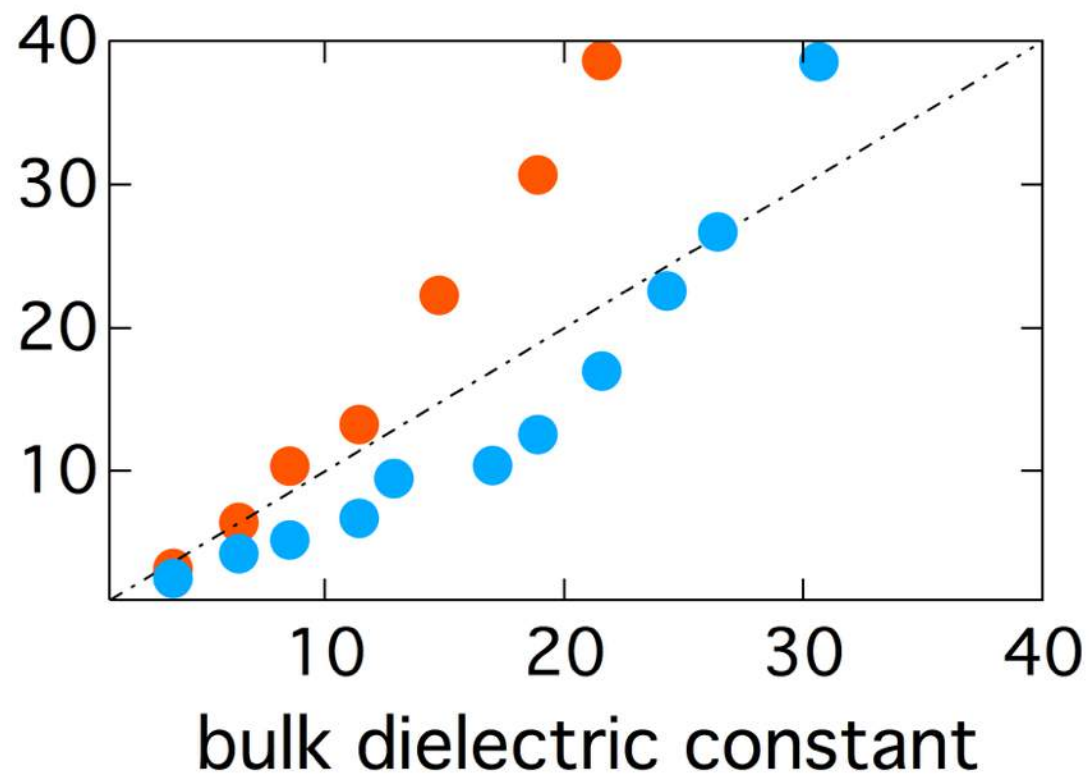
“Surface” dielectric constant

$$E^C(\mathbf{r}) = \sum_i q_i \underbrace{\Phi_s(\mathbf{r})}_{\text{solvent potential}}$$

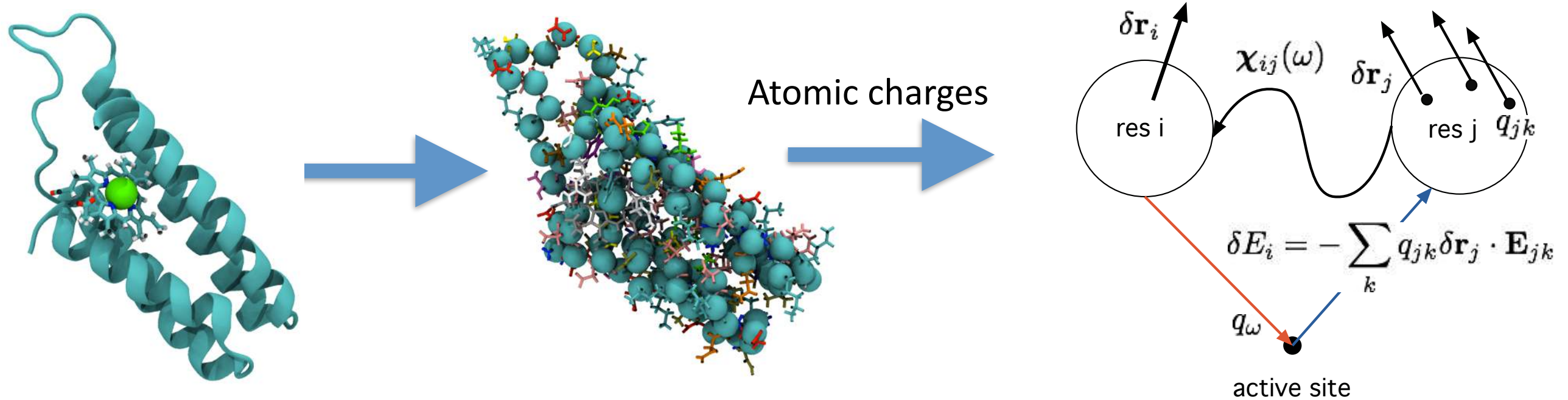
$$\sigma_P = \frac{1}{k_B T} \langle \delta P_n \delta E^C \rangle$$



surface diel. conts.



Dissipative **E**lectro-Elastic **N**etwork **M**odel (DENM)



Spring Constant

$$\delta r = r - r_0$$

$$E = \frac{C}{2} \sum_{i,j} H_{ij}^{\alpha\beta} \delta r_i^\alpha \delta r_j^\beta$$

Hessian Matrix

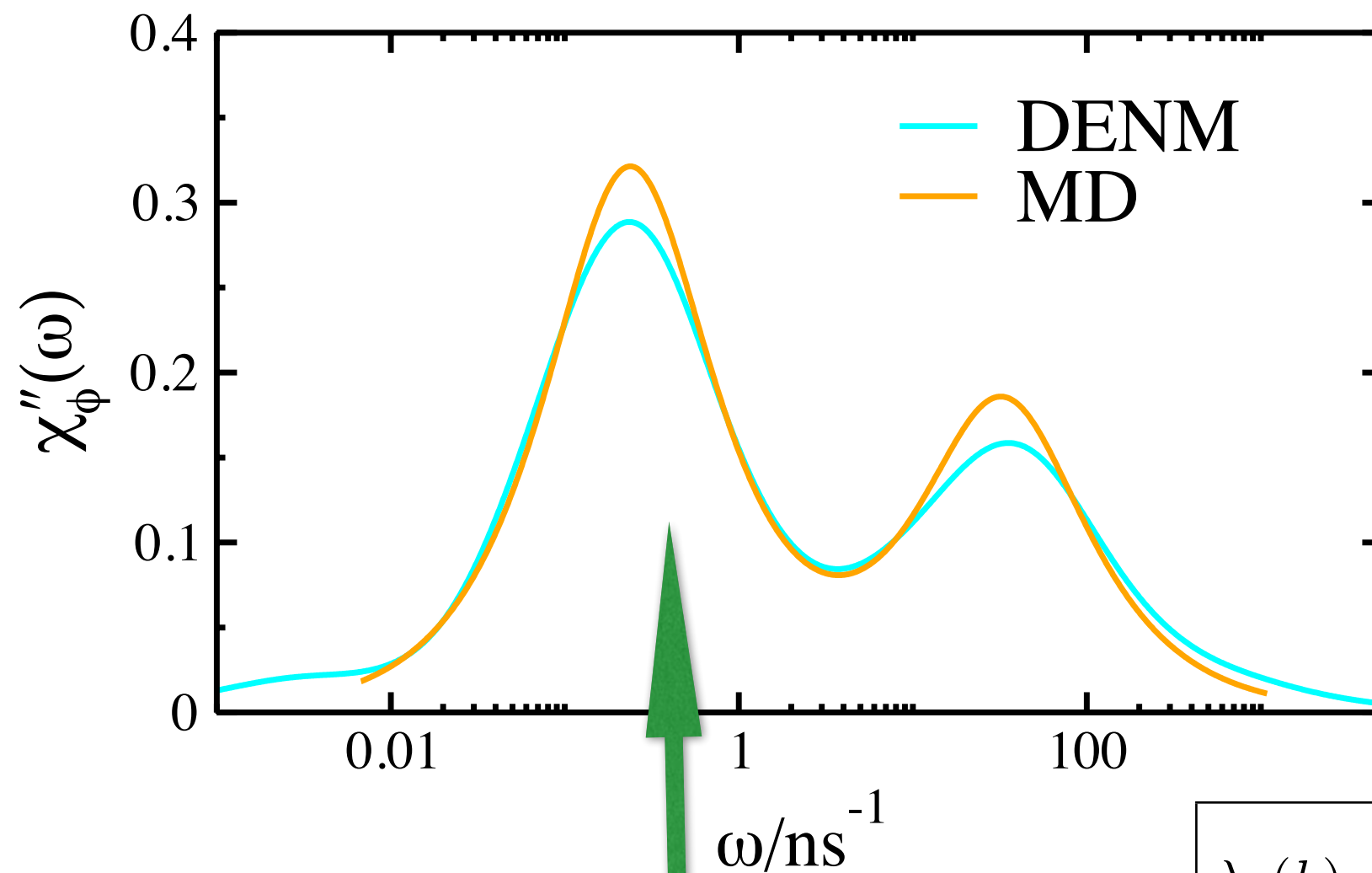
$$\int_0^t \zeta(t-t') \dot{\mathbf{q}}_m(t') dt' + \lambda_m \mathbf{q}_m = \mathbf{F}(t) + \mathbf{R}(t)$$

$$\chi_{ij}^{\alpha\beta}(\omega) = C^{-1} \sum_m U_{mi}^{\gamma\alpha} [\lambda_m + i\omega\zeta(\omega)]^{-1} U_{mj}^{\gamma\beta}$$

$$\chi_\phi(\omega) = - \sum_{i,j} E_{0j}^\alpha \chi_{ij}^{\alpha\beta}(\omega) E_{0i}^\beta$$

CytB: potential response

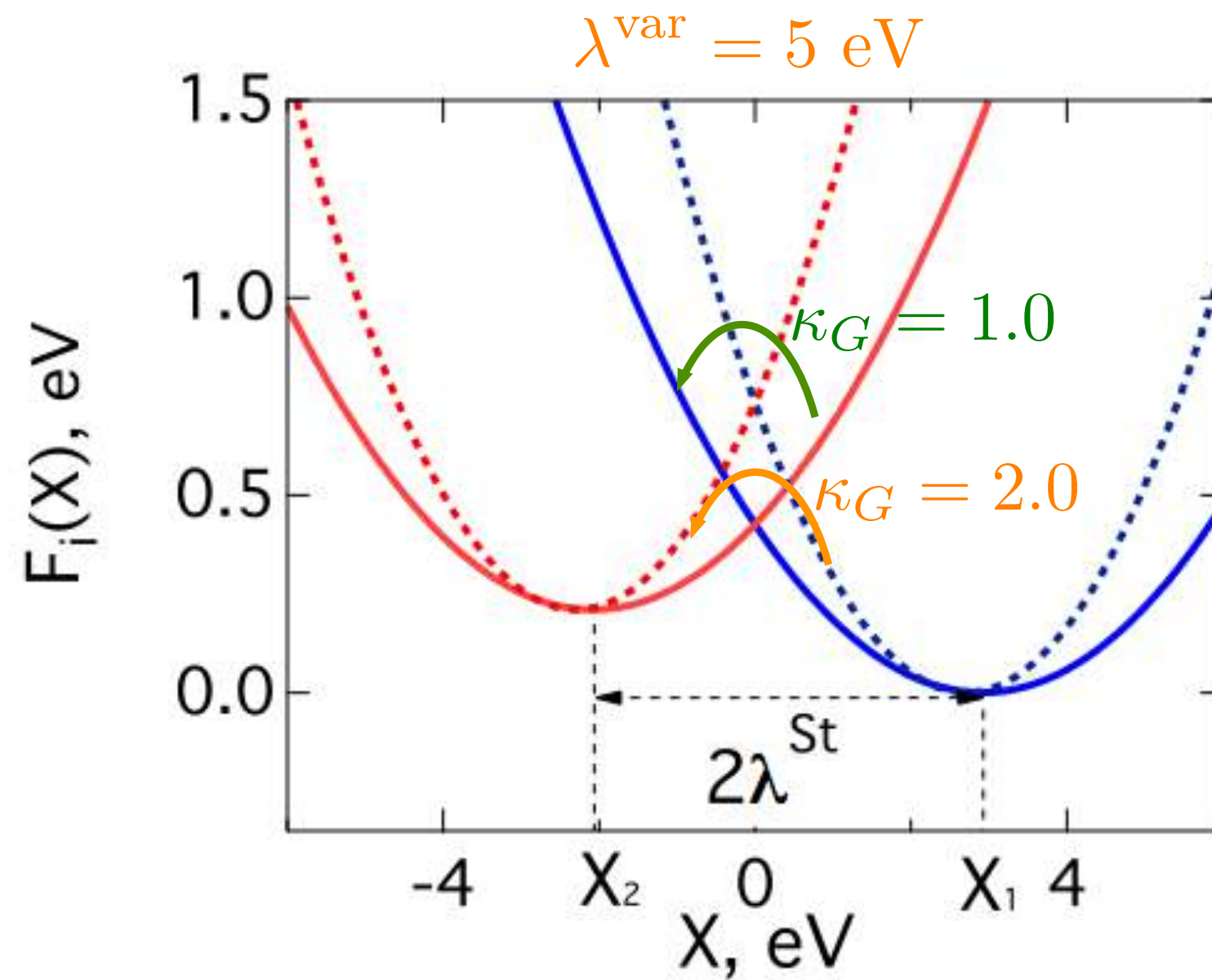
$$\chi_{\phi}(\omega) = - \sum_{i,j} E_{0j}^{\alpha} \chi_{ij}^{\alpha\beta}(\omega) E_{0i}^{\beta}$$



$$\lambda_s(k) \propto \int_{k=\omega_{\text{obs}}}^{\infty} \chi''(\omega) (d\omega/\omega)$$

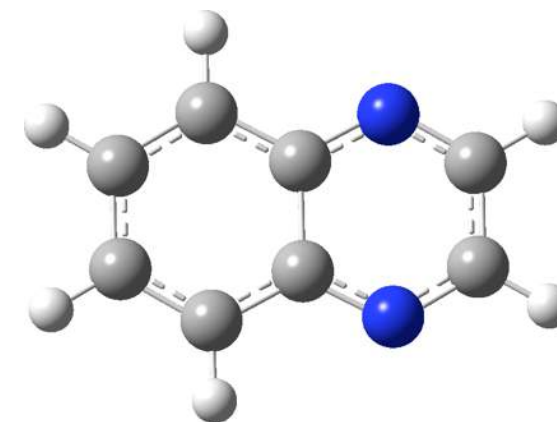
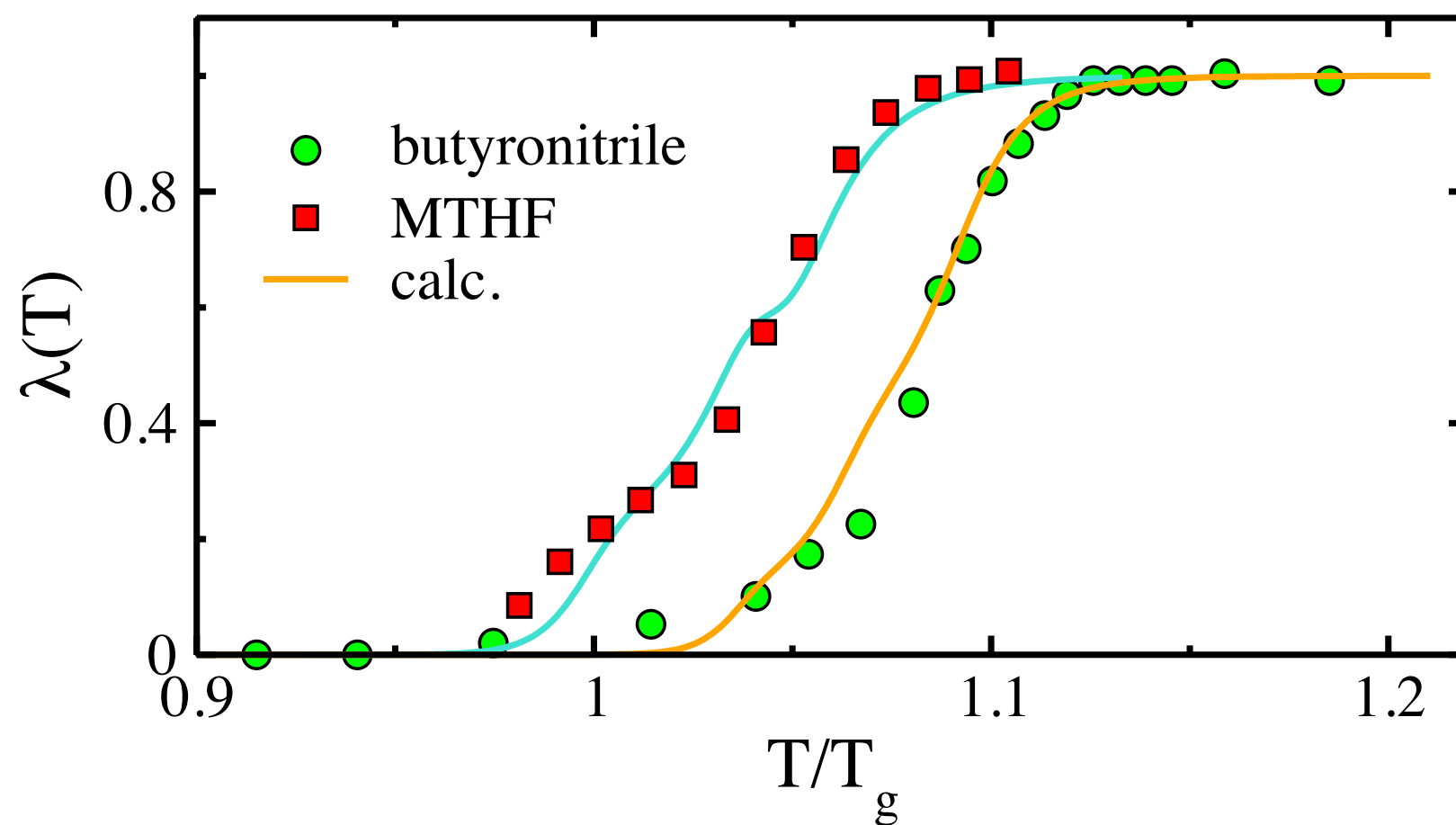
Elastic deformations of the protein shape

Nonergodic free energy surfaces



Spectroscopy in super-cooled liquids

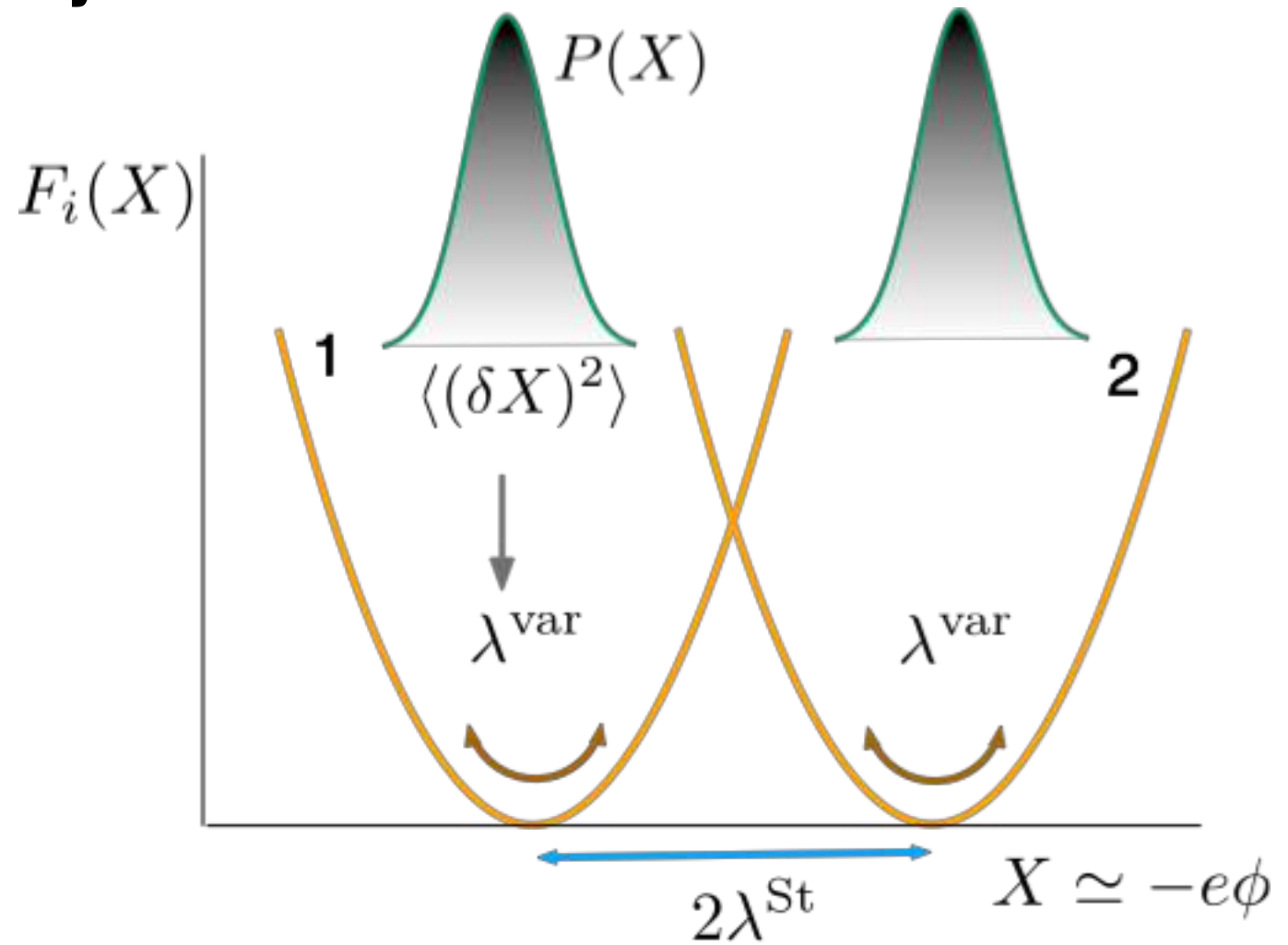
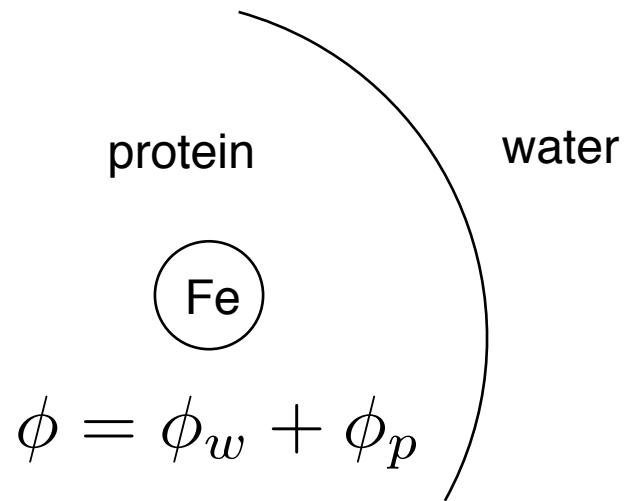
$$\omega_{\text{obs}} = 1/\tau_{\text{em}}$$



quinoxaline dye,
R. Richert, ASU

Freezing out of nuclear degrees of freedom on the time-scale
of phosphorescence

Signature of nonergodicity



$$\lambda^{\text{var}} = \beta \langle (\delta X)^2 \rangle / 2$$

$$2\lambda^{\text{St}} = X_1 - X_2$$

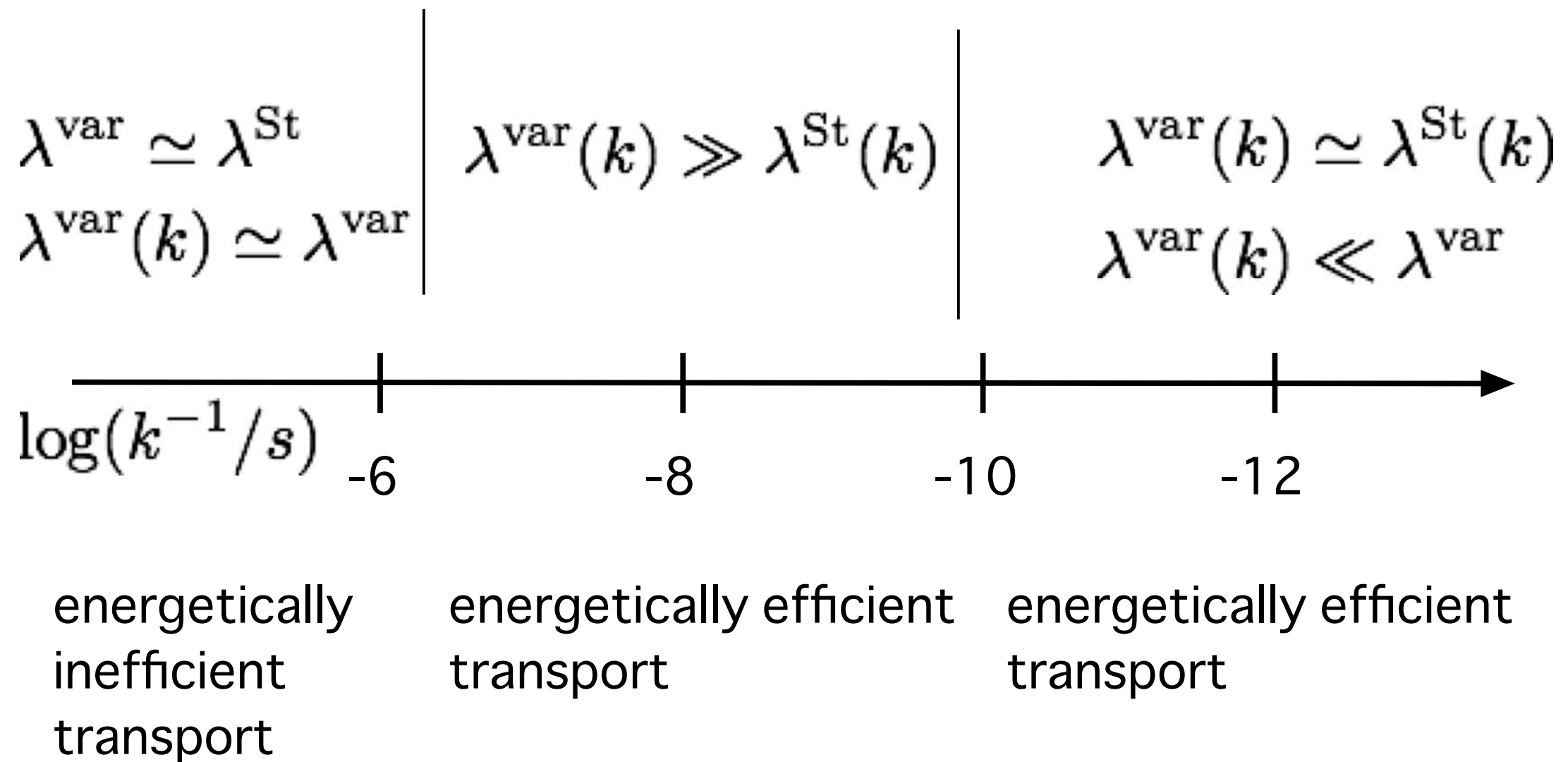
Canonical ensemble: $\lambda^{\text{St}} = \lambda^{\text{var}}$

$$\chi_G = \frac{\lambda^{\text{var}}}{\lambda^{\text{St}}}$$

$$\chi_G \gg 1$$

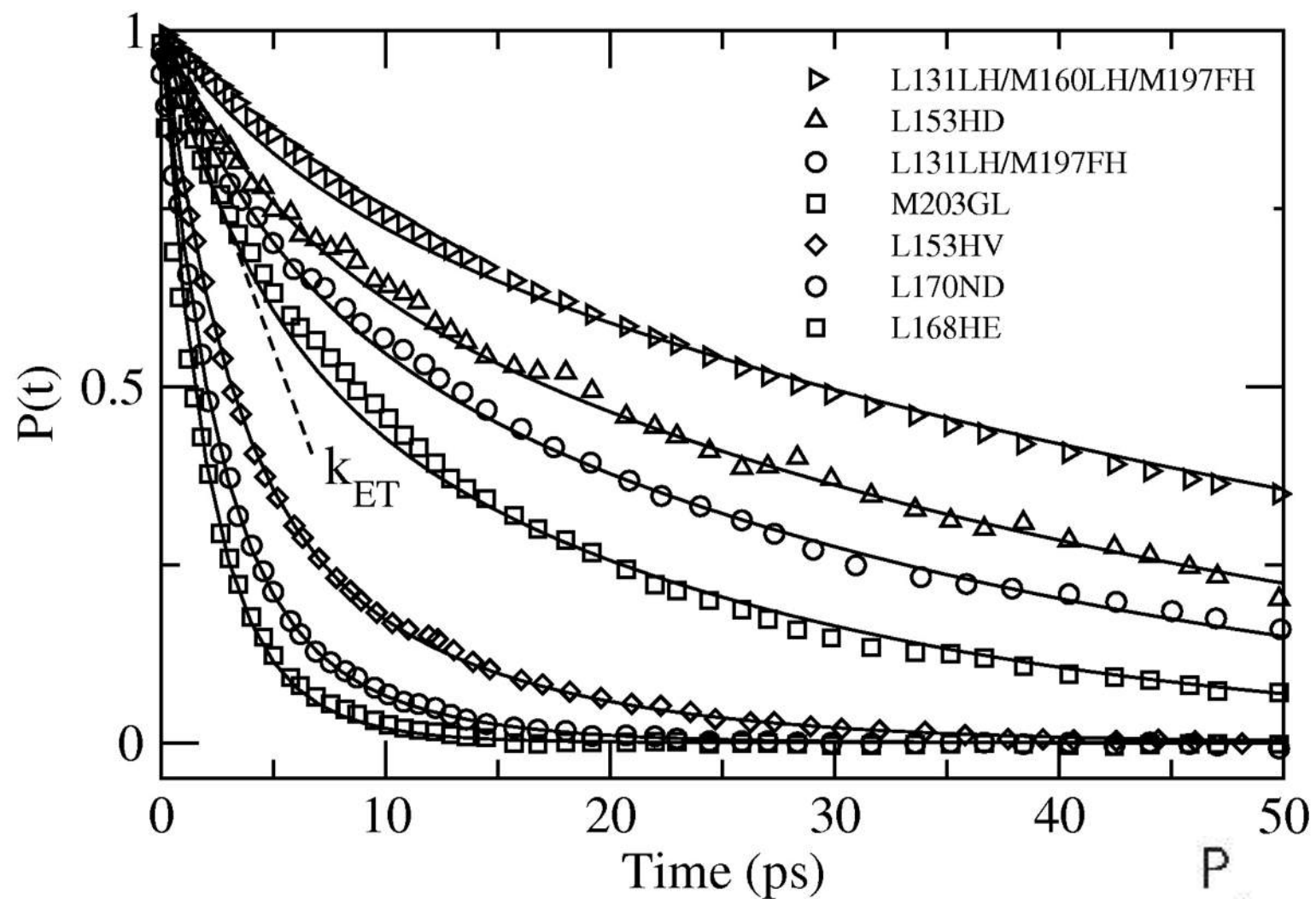
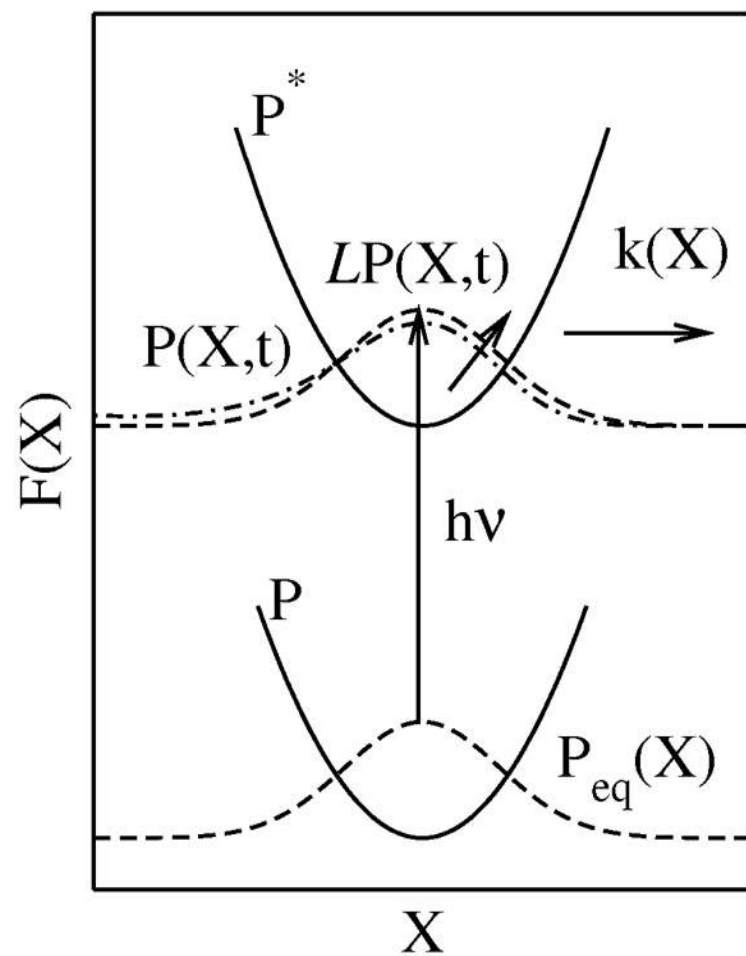
proteins

Time arrow of biological electron transport



$$\Delta G \rightarrow \Delta G / \kappa_G$$

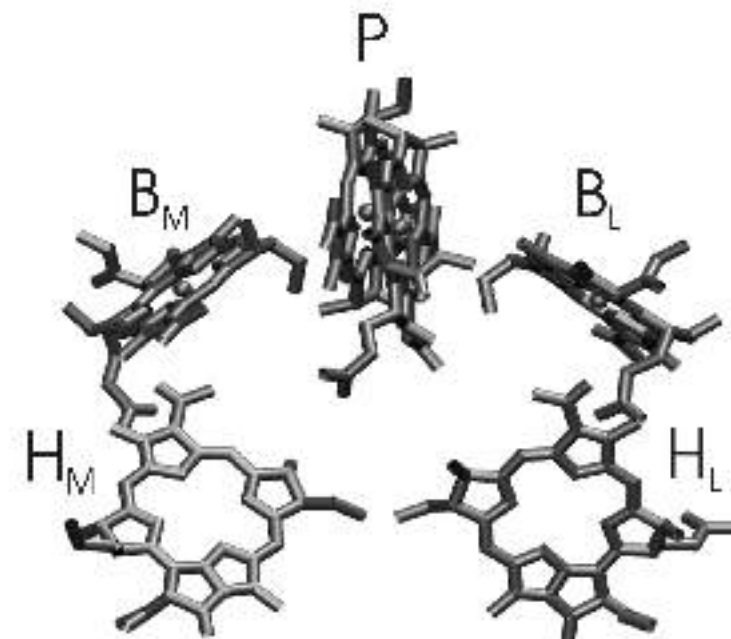
Population dynamics



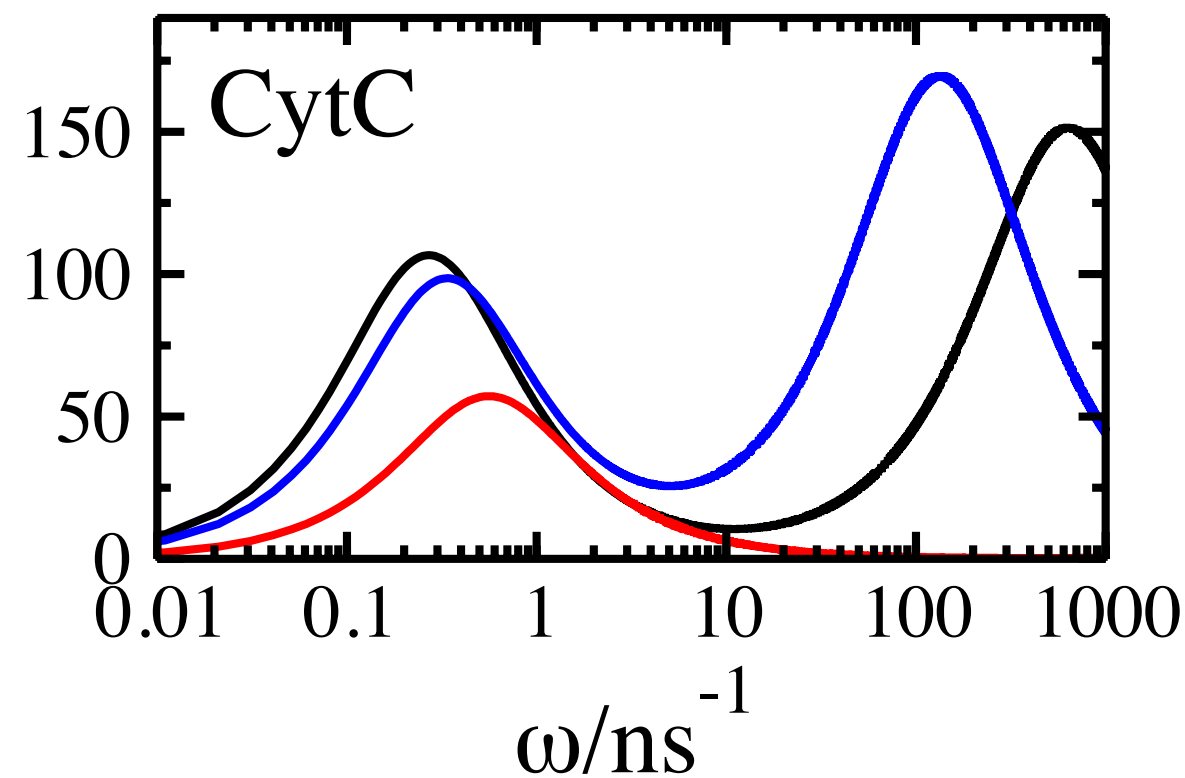
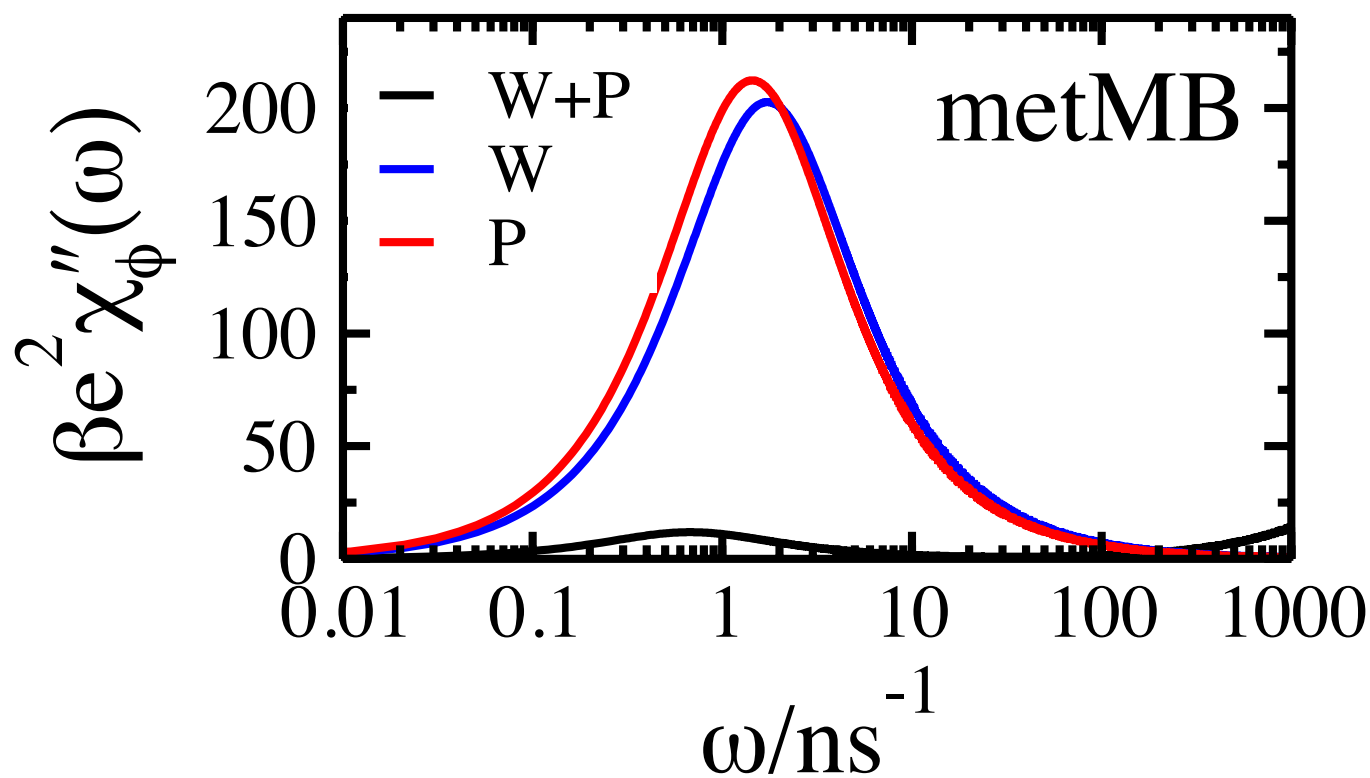
$$\partial P(X,t)/\partial t = [L(k_{\text{ET}}, X) - k(X)] P(X,t)$$

Fokker-Planck operator depending on $F(X, k_{\text{ET}})$

$$P(t) = \int P(X,t) dX$$



Does sequence matter?

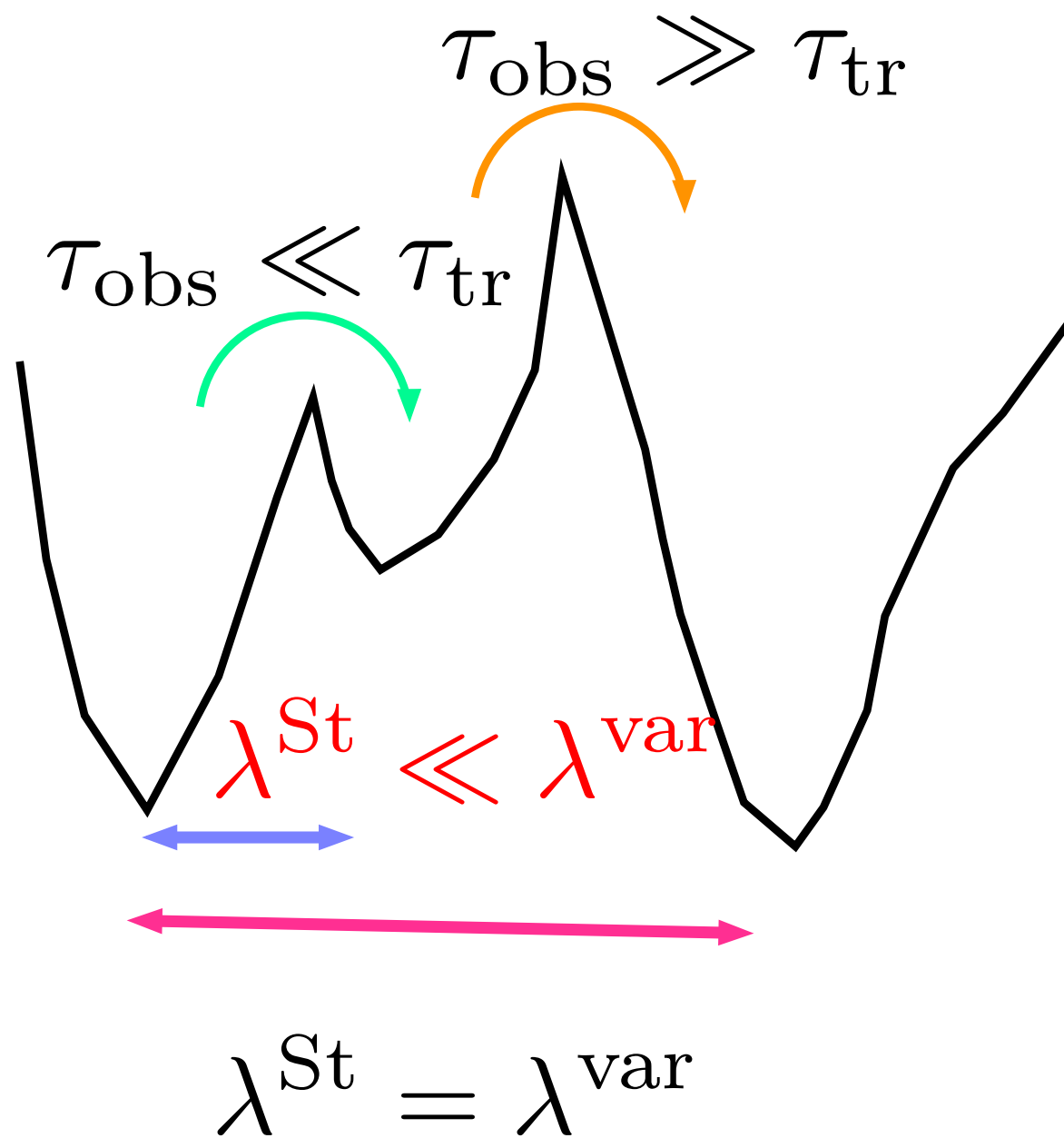


$$\lambda_s = \lambda_w + \lambda_p + \lambda_{pw} \propto \langle (\delta X_w + \delta X_p)^2 \rangle$$

Protein	λ_p, eV	λ_w, eV	λ_s, eV
Myoglobin	9.9	7.9	1.8
Cytochrome c	6.9	3.3	6.7

critical term

Protein-water compensation depends on the **surface charge** (sequence+folding).



Energy conversion machinery

